



Laser-Induced Current Transients in Strained-Si Diodes

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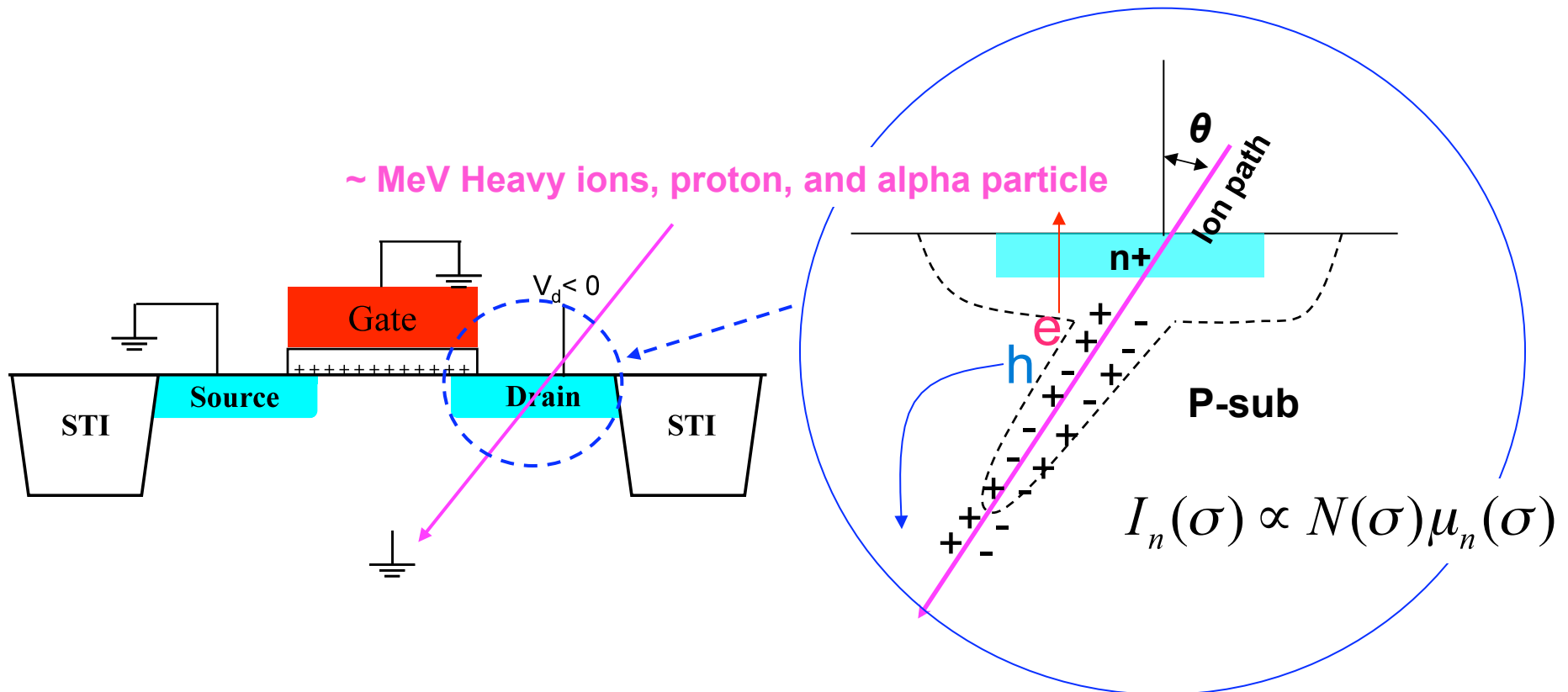
- Introduction
- Laser-induced current transients in diodes under mechanical stress
 - $\langle 110 \rangle$ uniaxial stress : Experiment and Simulation Results
 - $\langle 100 \rangle$ uniaxial and biaxial stress : Analytical Modeling
- Conclusion and future work

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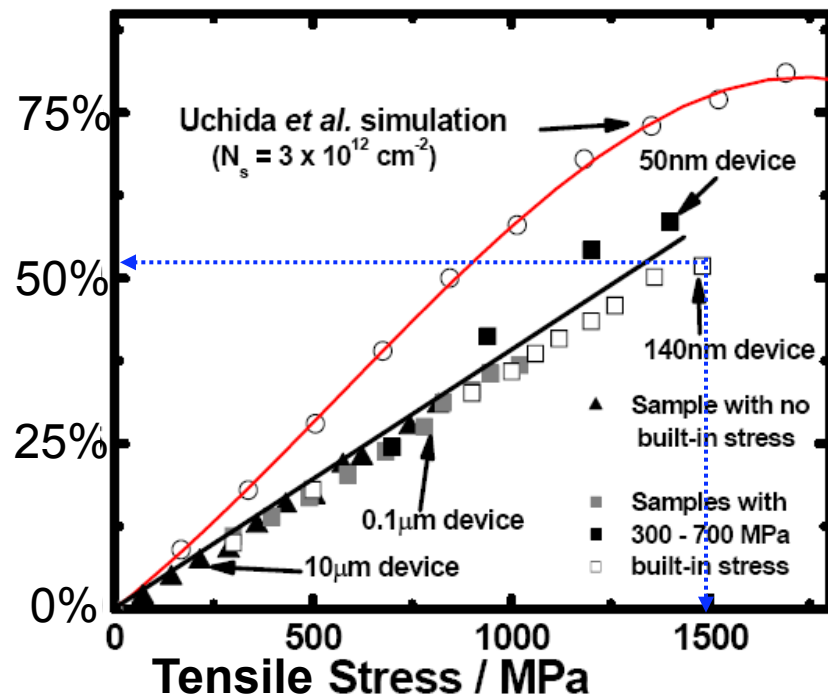
Single Event Effect in MOSFET



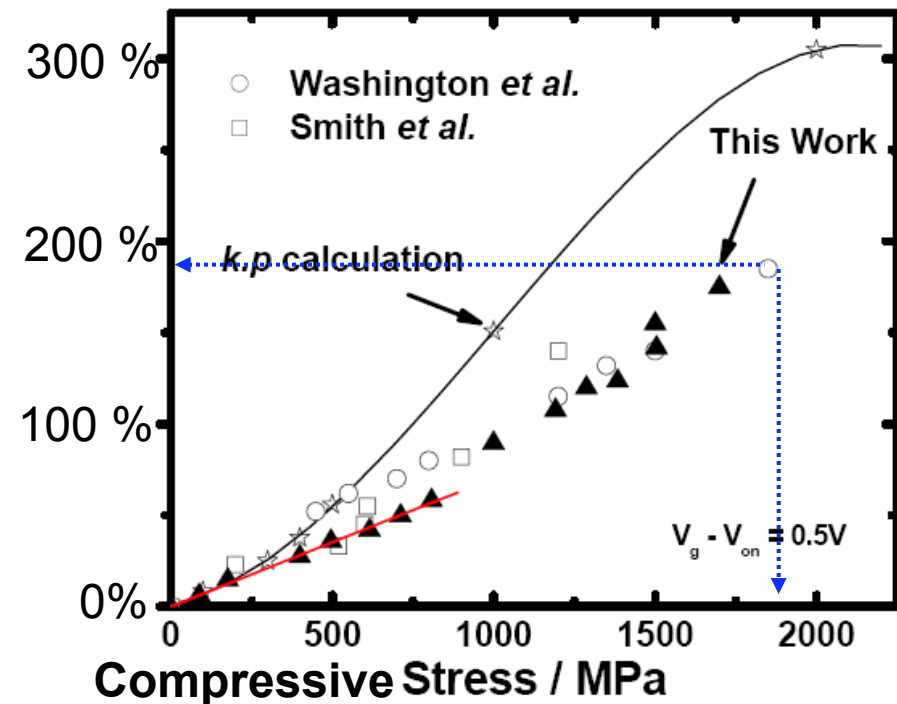
- Single Event Effect (transient pulse) occurs in MOSFET
- : Single Event Transient (SET) and Single Event Upset (SEU)
- Modern devices are sensitive to these effects.

Mechanical Stress Alters Mobility Significantly

Electron Mobility Enhancement



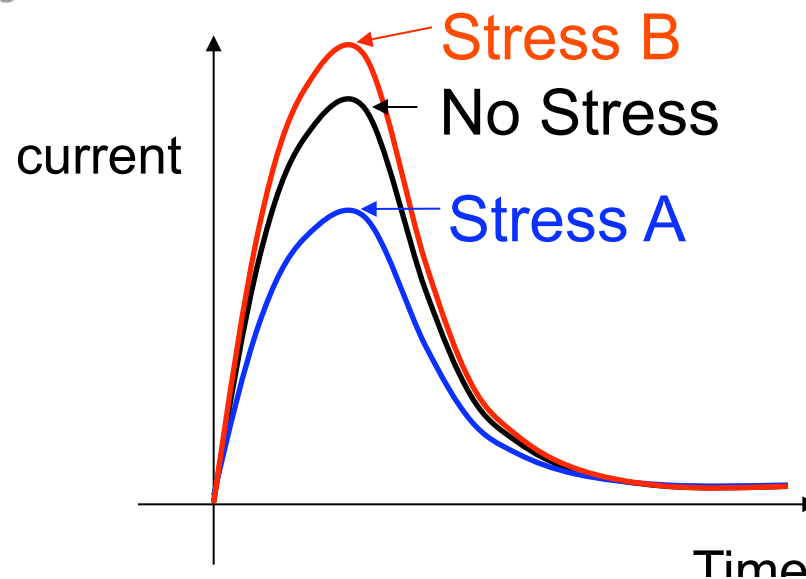
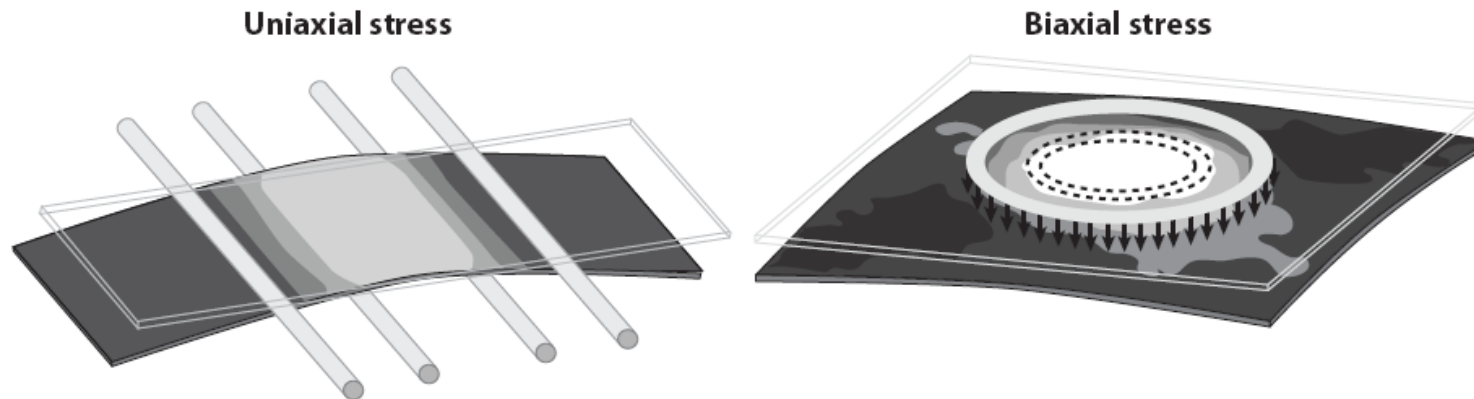
Hole Mobility Enhancement



Suthram, 2008

- In modern devices, the strained-Si technology is implemented to boost transistor performance.
- Mechanical stress enhances electron and hole mobility significantly (Performance Booster).
- What if we put this technologies in radiation environment?

Current Transients under Different Type of Stress



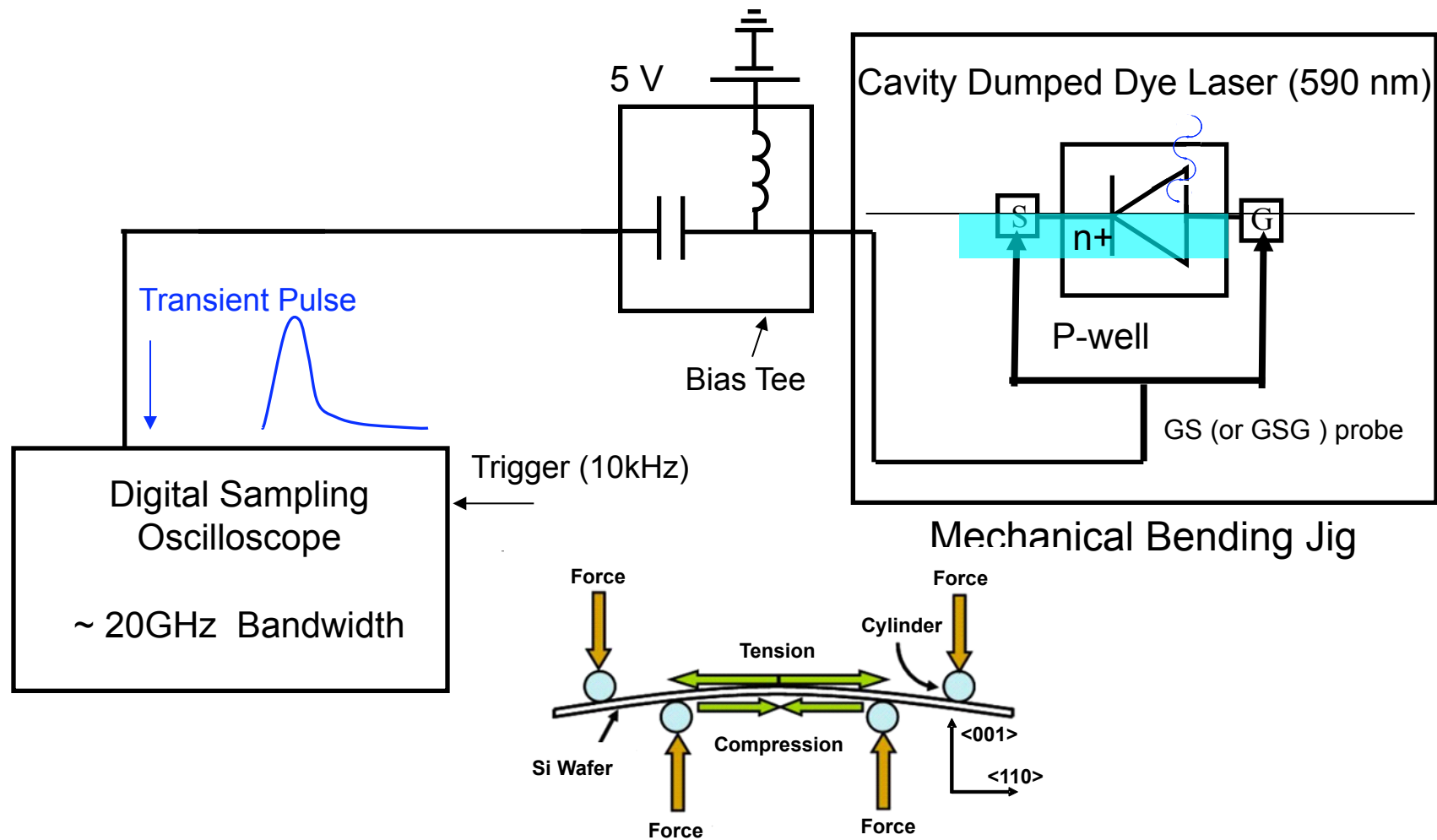
- How does different type of stress change current transient in diodes?
(I_{max} : peak current, Q: charge collection)

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Laser-Induced Current Transient Measurement System



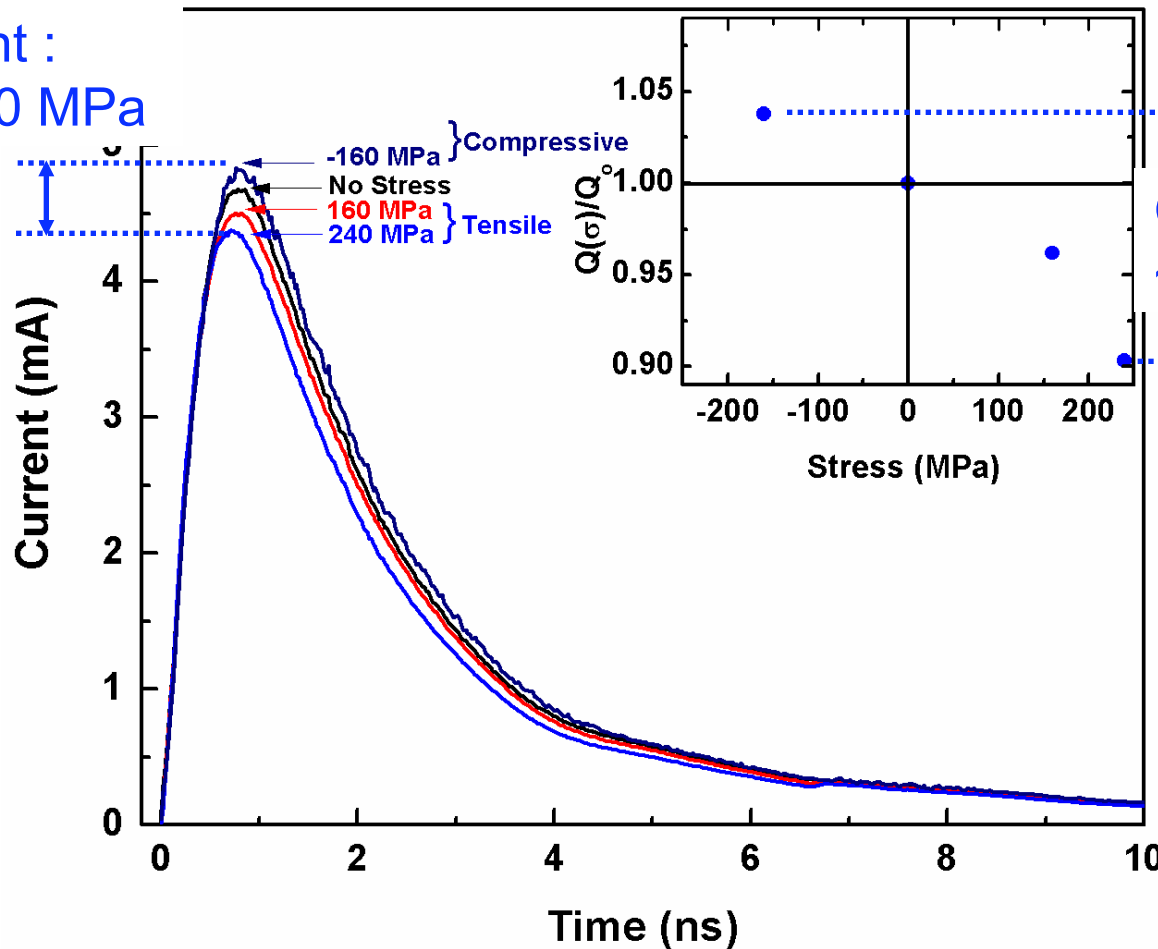
- We can observe changes of current transient pulses under mechanical stresses
- It is possible to apply different tensile or compressive stress using mechanical bending jig



Laser-induced Current Transients (Experiment)



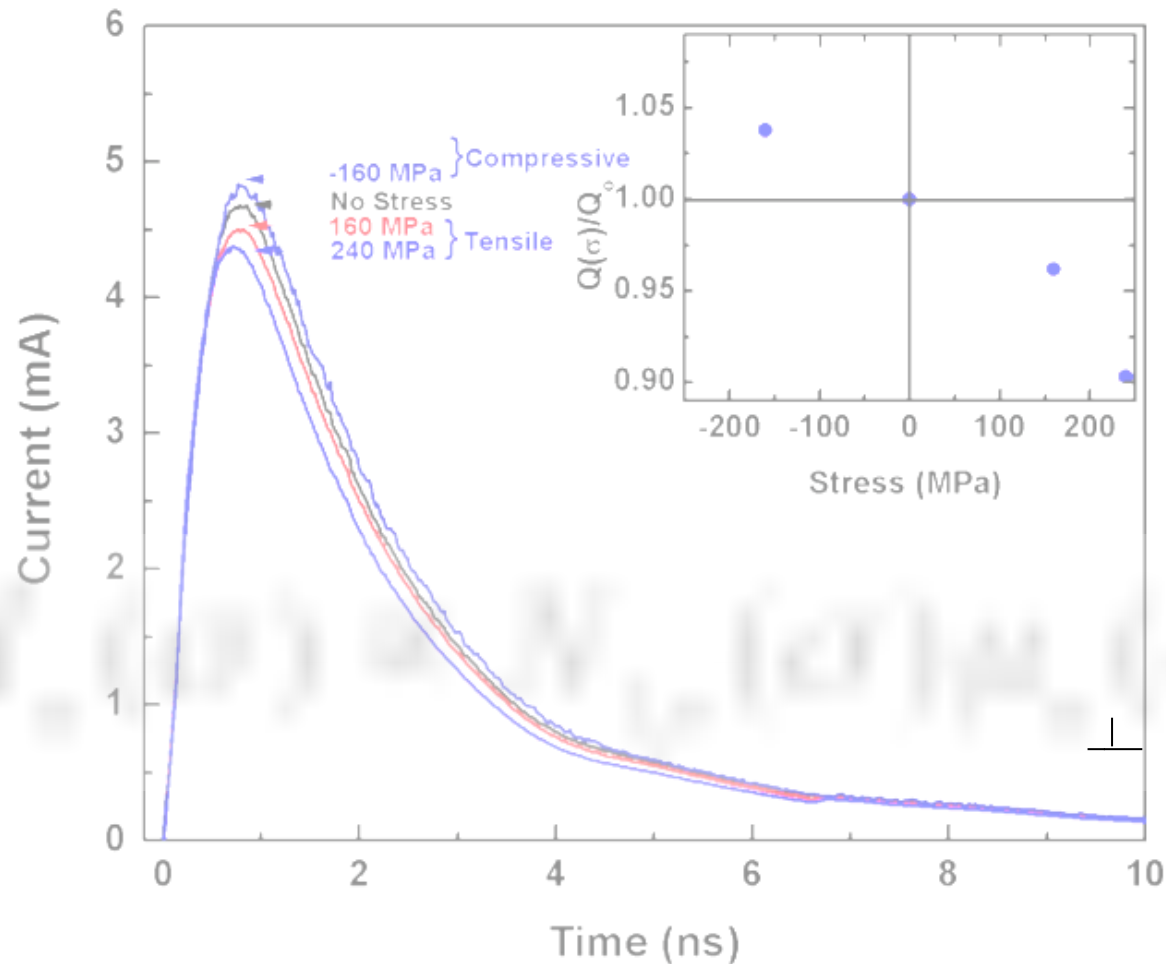
Peak current :
~11% at 400 MPa



Collected charges:
~14% at 400 MPa

- Tensile (Compressive) stress decrease (increase) peak current and collected charges.

Simplified Current



- $N_{1p}(\sigma)$ is related to bandgap narrowing under mechanical stress.
- Mechanical stress alters electron mobility along $\langle 001 \rangle$ direction.

Strain Effect on Electron-hole Pair Generation (Bandgap Narrowing)

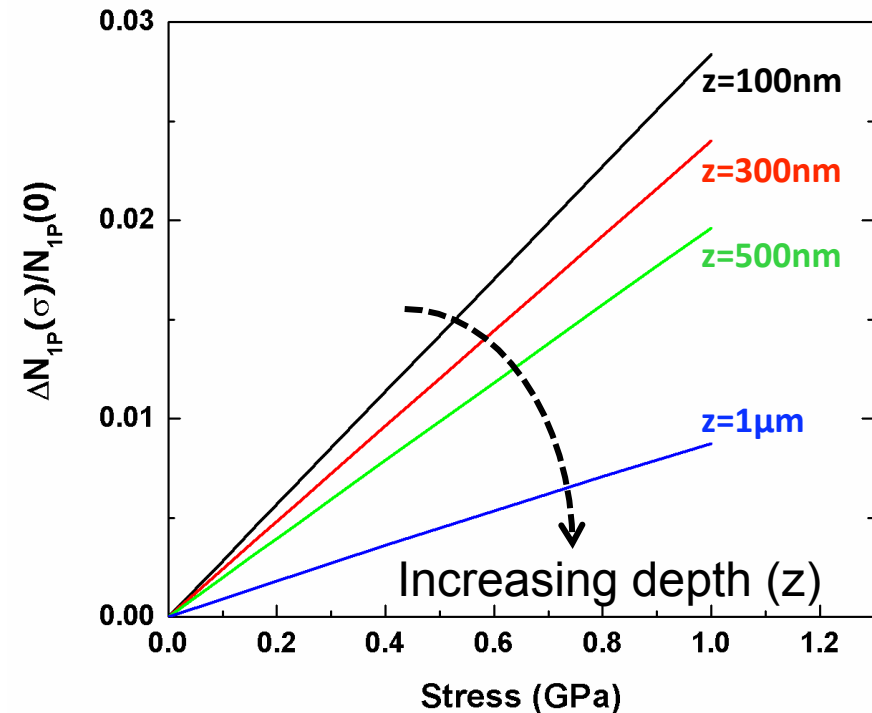
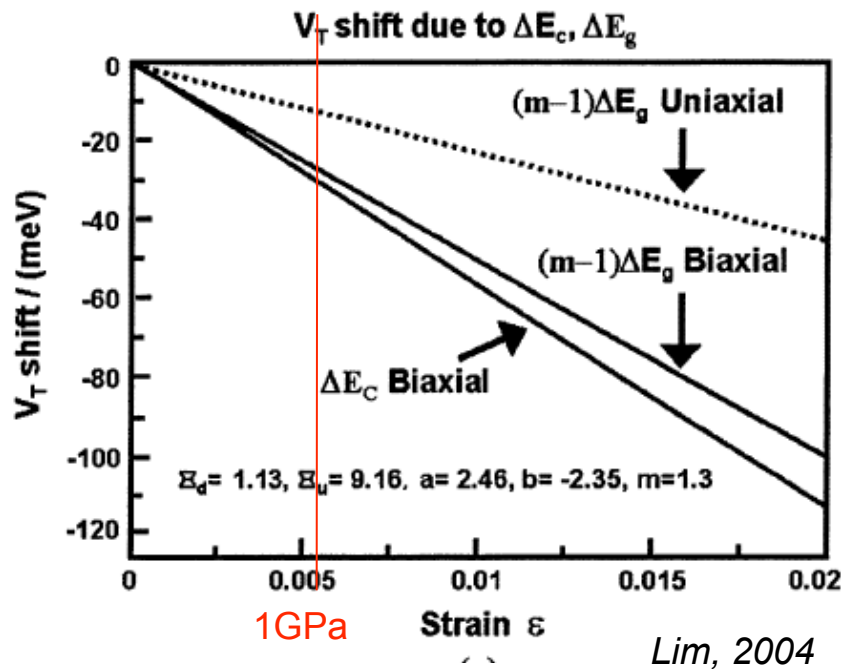


Absorption coefficient

$$\frac{\Delta\alpha}{\alpha} = \frac{\Delta E_g(\sigma)}{h\nu - E_g} \ll 1, \quad h\nu > E_g$$

The number of generated e-h pairs

$$N_{1p}(z) = \frac{\alpha}{\hbar\omega} \exp(-\alpha z) \int_{-\infty}^{\infty} I_0(z, t) dt$$

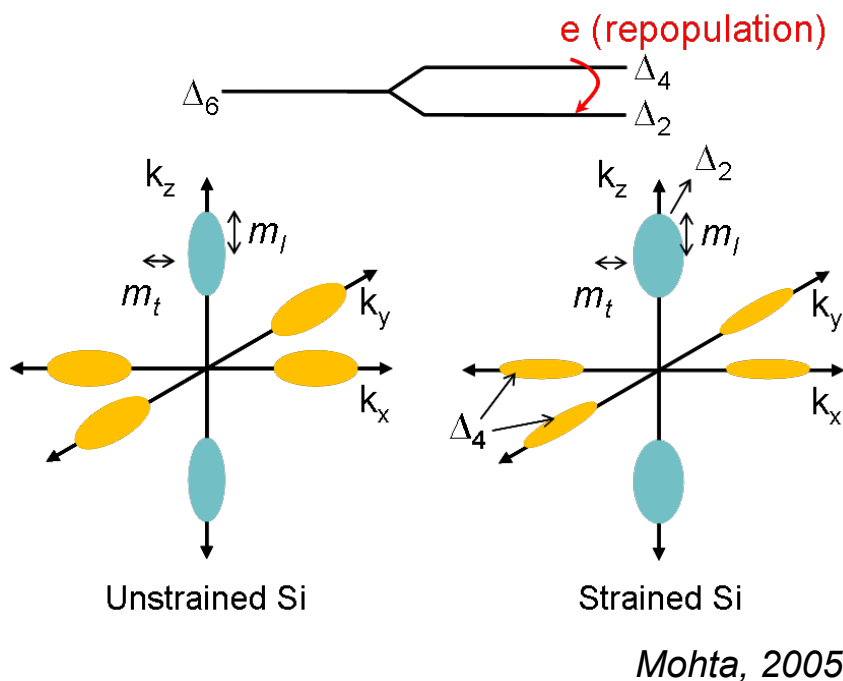


$\Delta E_g \cong 30 \text{ meV}$ at 1 GPa of uniaxial tensile stress

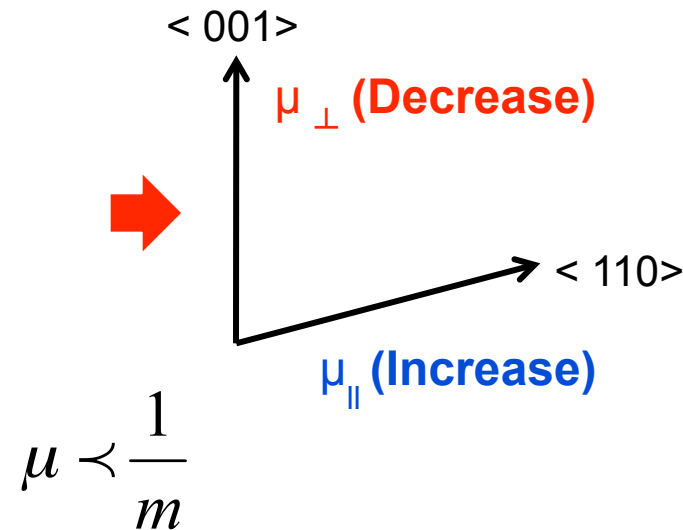
- Change in N_{1p} is less than 3% for 1 GPa of uniaxial tensile stress

Strain Effect on Electron Mobility

Under Tensile Stress :



$$I_n(\sigma) \propto N_{1p}(\sigma) \mu_{n\perp}(\sigma)$$

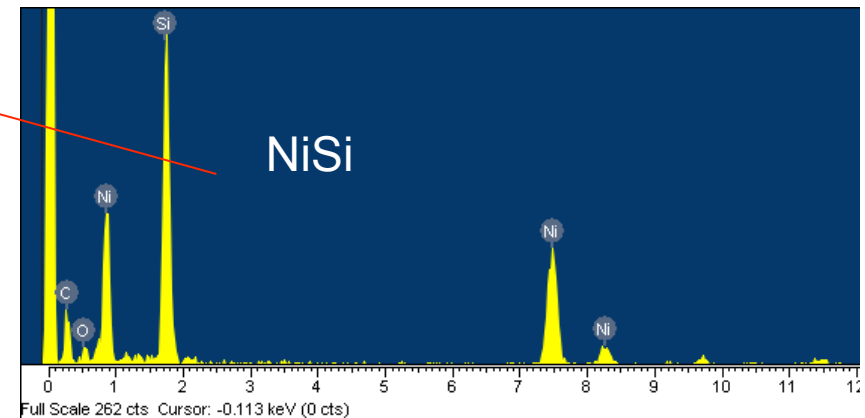
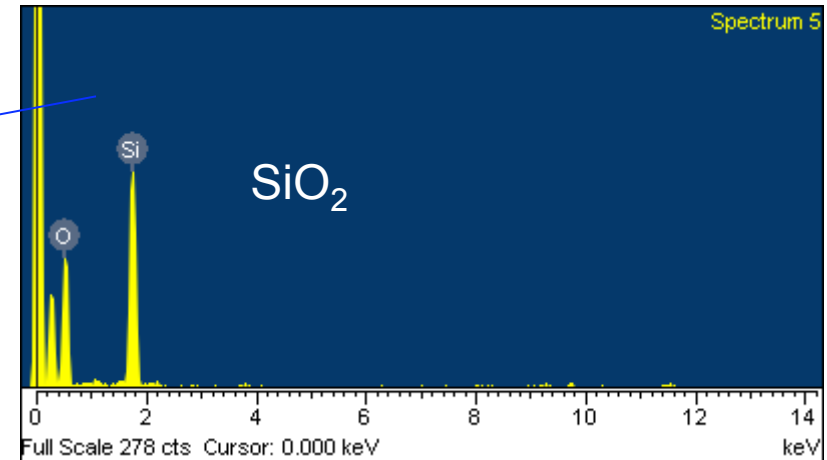
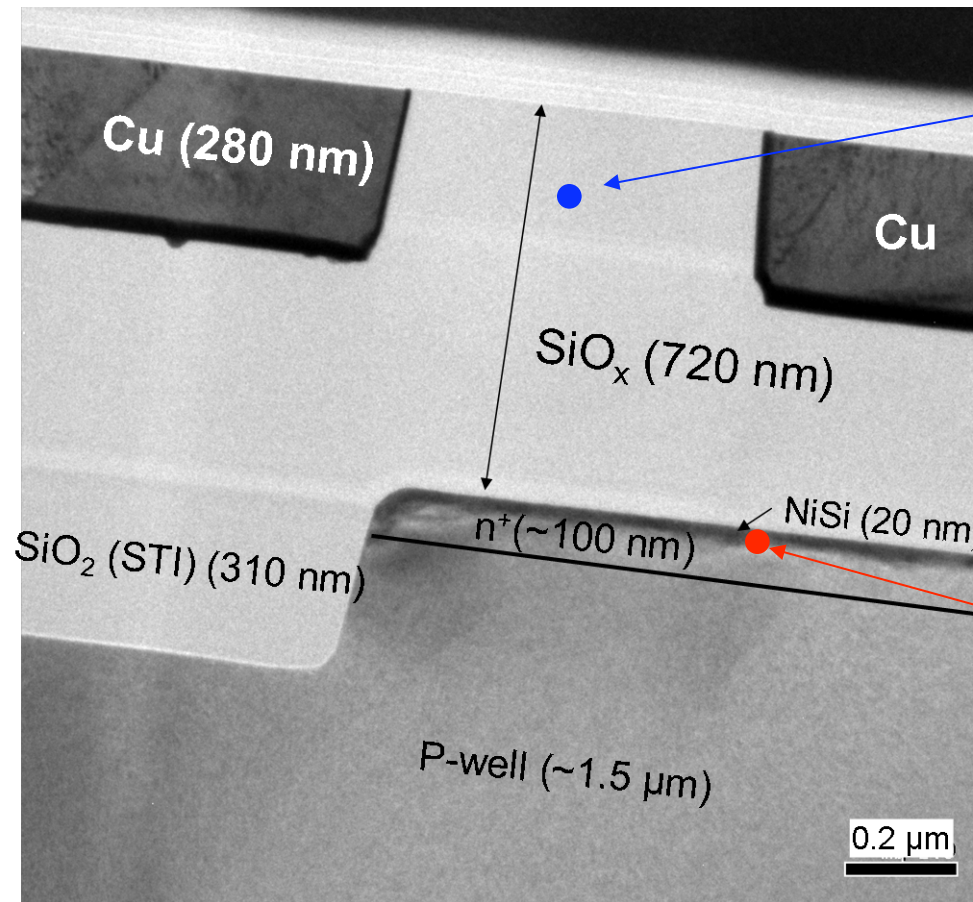


For 1 GPa of uniaxial tensile stress,

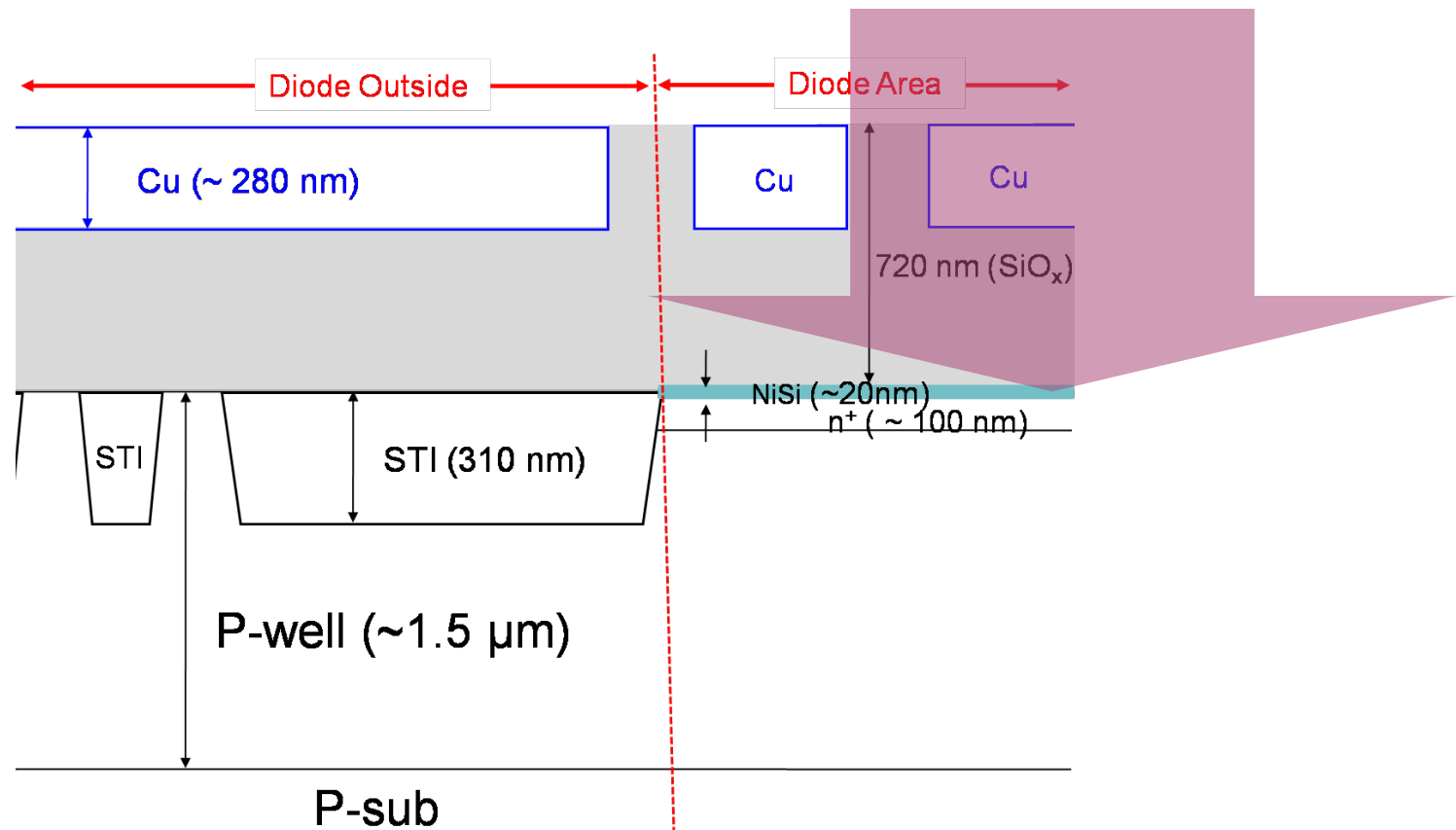
- electron mobility changes ~53% and Number of e-h pairs < 3%

→ Electron mobility change dominates current transient change under stress

Diode Structure through TEM and EDS

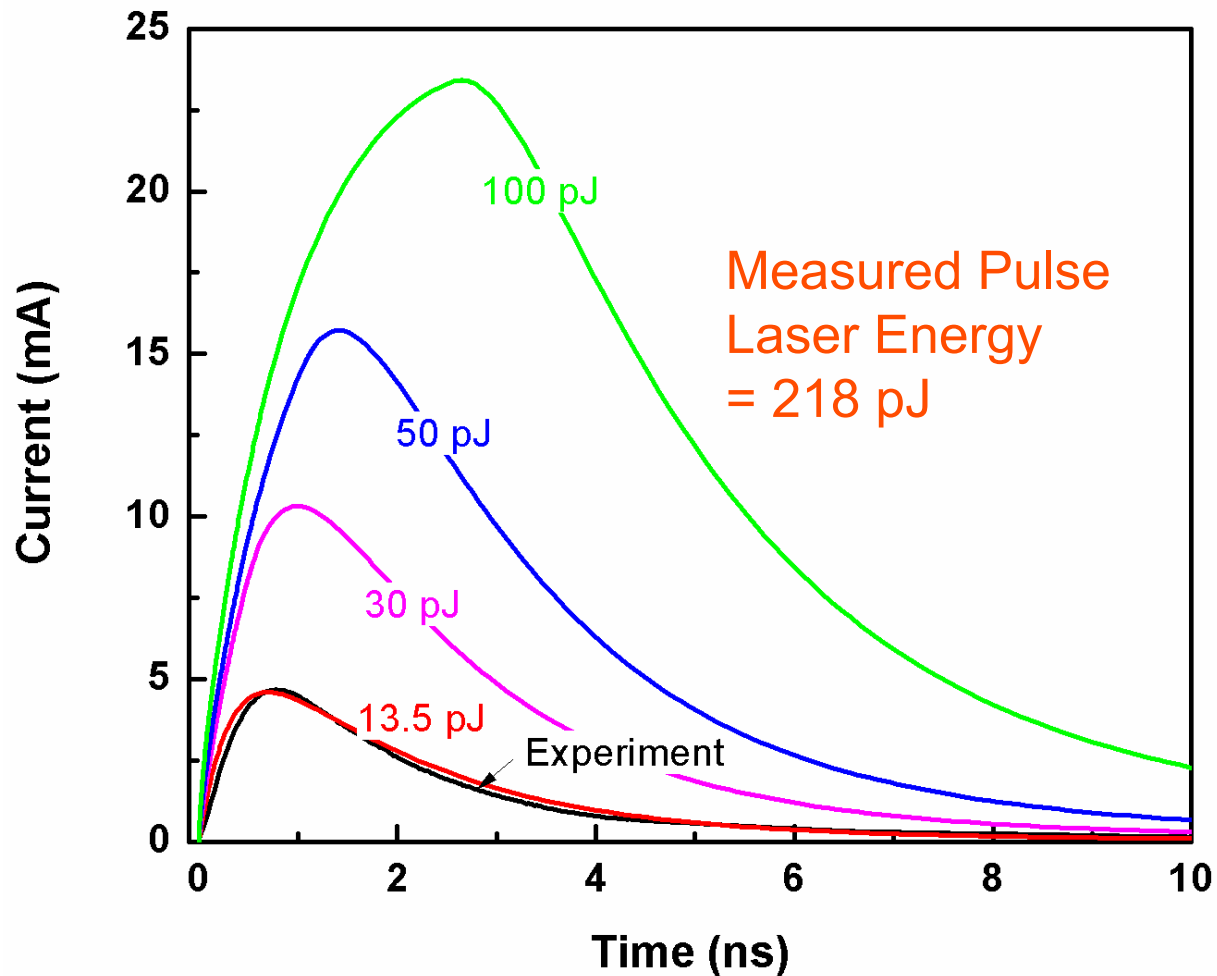


Deposition of Pulse Laser into n⁺p diode



- Some of pulse laser energy are reflected and absorbed in layers on top of a diode.

Current Transients according to Pulse Laser Energy



- The discrepancy between simulated energy and measured energy is shown.

Laser Pulse Energy into Si



$$\begin{aligned} I_{Si} &= I_{\text{intial}} \times (\text{transmission due to Cu/SiO}_2) \\ &\quad \times (\text{transmission in SiO}_2) \times (\text{transmission in NiSi}) \\ &= I_{\text{intial}} \times (0.57) \times (0.80) \times (0.51 \times 0.23) \\ &= 0.056 I_{\text{intial}} \end{aligned}$$

Intensity (I) \propto Power (P) \propto Laser Pusle Energy (E)

Measured Laser Pulse Energy = 216 pJ

Pulse Laser Energy deposited in Si = (216 pJ) \times (0.054)

= 12 pJ

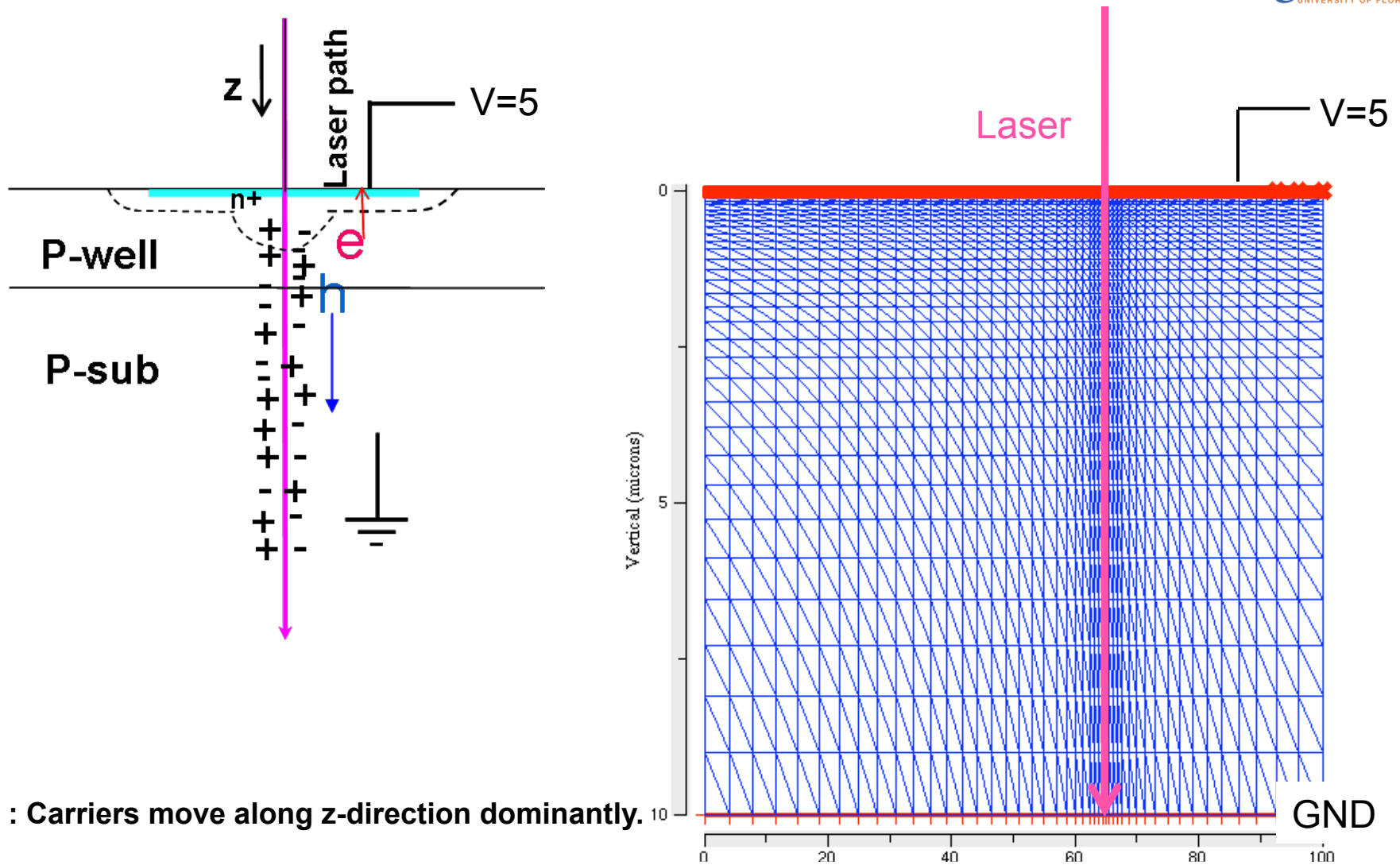
Melinger, 1994
Amiotti, 1990

If we consider uncertainty of every layer characteristics, $E = 10 \sim 25$ pJ

Especially, consider the optical characteristics nickel silicide (NiSi, NiSi₂, Ni₂Si)

2D Simulation value => 13.5 pJ (It is very close to calculated values above)

2-D Simulation Structure of n⁺p Diode



*Note : Carriers move along z-direction dominantly.

- Using analysis from experimental results such as piezoresistance coefficient and bandgap narrowing effect, we did 2-D FLOOD current transient simulation under stress.

Strain Model (Piezoresistance Effect)



$$\begin{bmatrix} \pi_{11} & \pi_{12} & \pi_{12} & 0 & 0 & 0 \\ \pi_{12} & \pi_{11} & \pi_{12} & 0 & 0 & 0 \\ \pi_{12} & \pi_{12} & \pi_{11} & 0 & 0 & 0 \\ 0 & 0 & 0 & \pi_{44} & 0 & 0 \\ 0 & 0 & 0 & 0 & \pi_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & \pi_{44} \end{bmatrix} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \sigma_{yz} \\ \sigma_{zx} \\ \sigma_{xy} \end{bmatrix} = \begin{bmatrix} \Delta\rho_{xx} / \rho_{xx} \\ \Delta\rho_{yy} / \rho_{yy} \\ \Delta\rho_{zz} / \rho_{zz} \\ \Delta\rho_{yz} / \rho_{yz} \\ \Delta\rho_{zx} / \rho_{zx} \\ \Delta\rho_{xy} / \rho_{xy} \end{bmatrix} = \begin{bmatrix} -\Delta\mu_{xx} / \mu_{xx} \\ -\Delta\mu_{yy} / \mu_{yy} \\ -\Delta\mu_{zz} / \mu_{zz} \\ -\Delta\mu_{yz} / \mu_{yz} \\ -\Delta\mu_{zx} / \mu_{zx} \\ -\Delta\mu_{xy} / \mu_{xy} \end{bmatrix}$$

$$\begin{bmatrix} \frac{\rho_{xx} + \Delta\rho_{xx}}{\rho_{xx}} & \frac{\Delta\rho_{xy}}{\rho_{xy}} & \frac{\Delta\rho_{xz}}{\rho_{xz}} \\ \frac{\Delta\rho_{xy}}{\rho_{xy}} & \frac{\rho_{yy} + \Delta\rho_{yy}}{\rho_{yy}} & \frac{\Delta\rho_{yz}}{\rho_{yz}} \\ \frac{\Delta\rho_{xz}}{\rho_{xz}} & \frac{\Delta\rho_{yz}}{\rho_{yz}} & \frac{\rho_{zz} + \Delta\rho_{zz}}{\rho_{zz}} \end{bmatrix} \begin{bmatrix} J_x(0) \\ J_y(0) \\ J_z(0) \end{bmatrix} = \begin{bmatrix} \frac{\mu_{xx} - \Delta\mu_{xx}}{\mu_{xx}} & \frac{-\Delta\mu_{xy}}{\mu_{xy}} & \frac{-\Delta\mu_{xz}}{\mu_{xz}} \\ \frac{-\Delta\mu_{xy}}{\mu_{xy}} & \frac{\mu_{yy} - \Delta\mu_{yy}}{\mu_{yy}} & \frac{-\Delta\mu_{yz}}{\mu_{yz}} \\ \frac{-\Delta\mu_{xz}}{\mu_{xz}} & \frac{-\Delta\mu_{yz}}{\mu_{yz}} & \frac{\mu_{zz} - \Delta\mu_{zz}}{\mu_{zz}} \end{bmatrix} \begin{bmatrix} J_x(0) \\ J_y(0) \\ J_z(0) \end{bmatrix} = \begin{bmatrix} J_x(\sigma) \\ J_y(\sigma) \\ J_z(\sigma) \end{bmatrix}$$

Mason, 1957
Smith, 1954

2-D Piezoresistance Coefficients



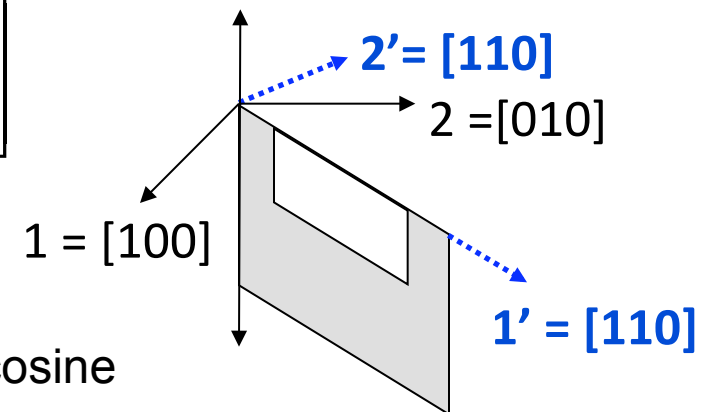
[100] Orientation “Smith” Coefficients:

Material	n-Si	p-Si
π_{11}	-102.2	6.6
π_{12}	53.4	-1.1
π_{44}	-13.6	138.1

10^{-5} MPa^{-1}

$$[\pi_{ij}] = \begin{bmatrix} \pi_{11} & \pi_{12} & 0 \\ \pi_{12} & \pi_{11} & 0 \\ 0 & 0 & \pi_{44} \end{bmatrix}$$

$3 = 3' = [001]$



Using directional cosine transformation

[110] Orientation Coefficients:

Material	n-Si	p-Si
π'_{11}	-31.2	71.8
π'_{13}	53.4	-1.1
π'_{33}	-102.2	6.6
π'_{55}	-13.6	138.1

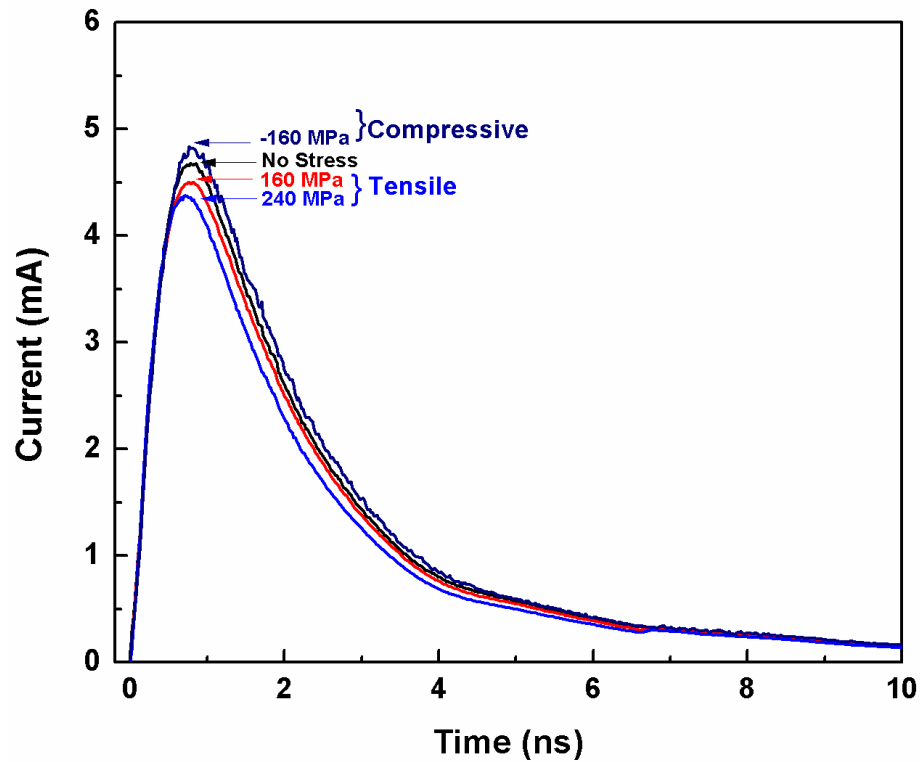
10^{-5} MPa^{-1}

$$[\pi'_{ij}] = \begin{bmatrix} \pi'_{11} & \pi'_{13} & 0 \\ \pi'_{13} & \pi'_{33} & 0 \\ 0 & 0 & \pi'_{55} \end{bmatrix}$$

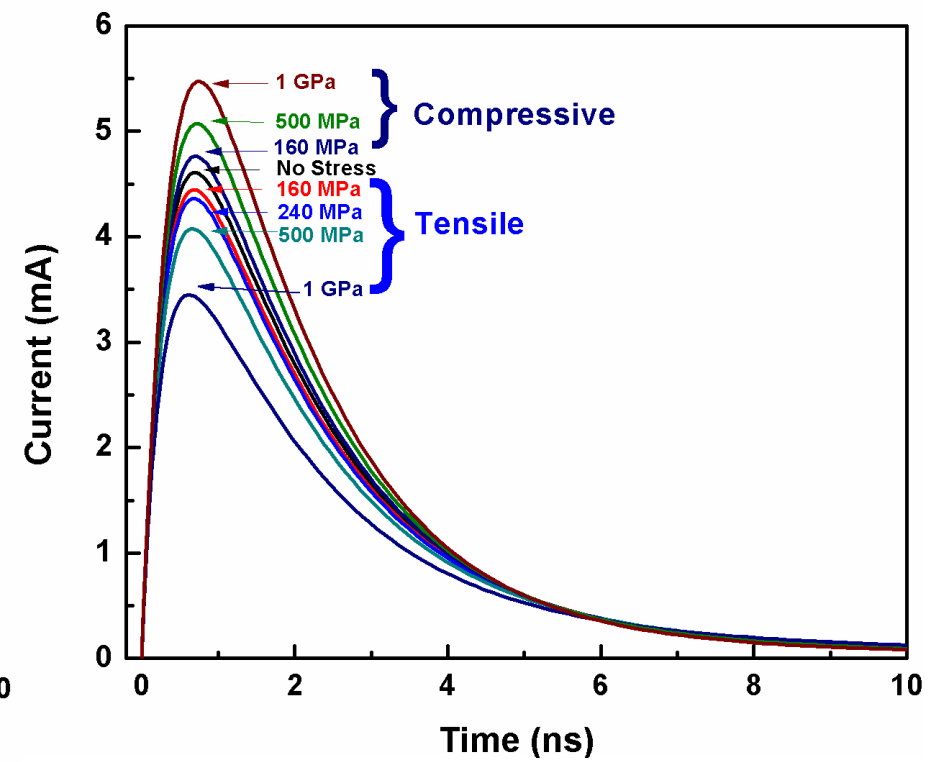
Smith, 1954
Kanda, 1982

Experiment vs. 2-D Simulation results

Experiment

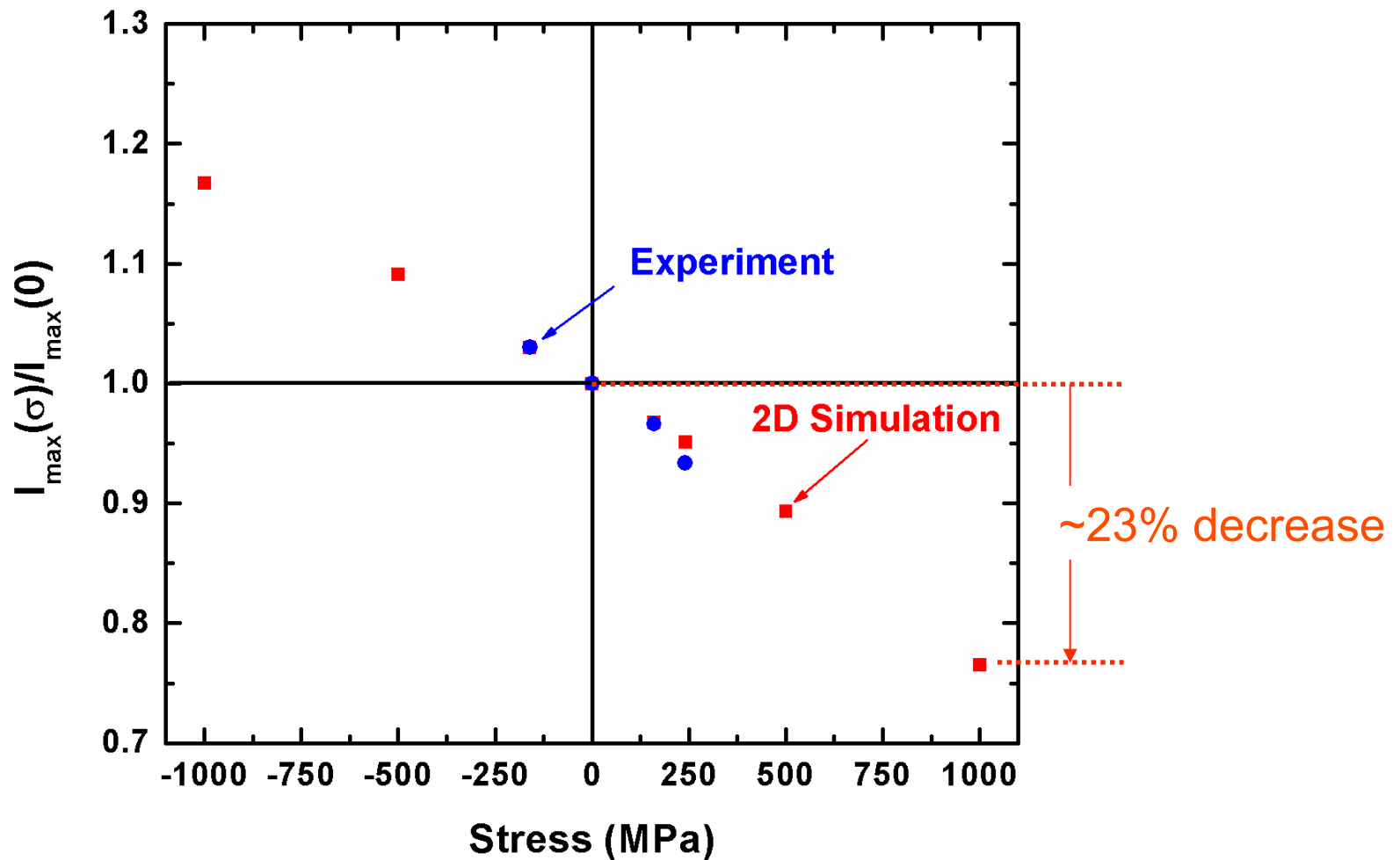


2-D Simulation



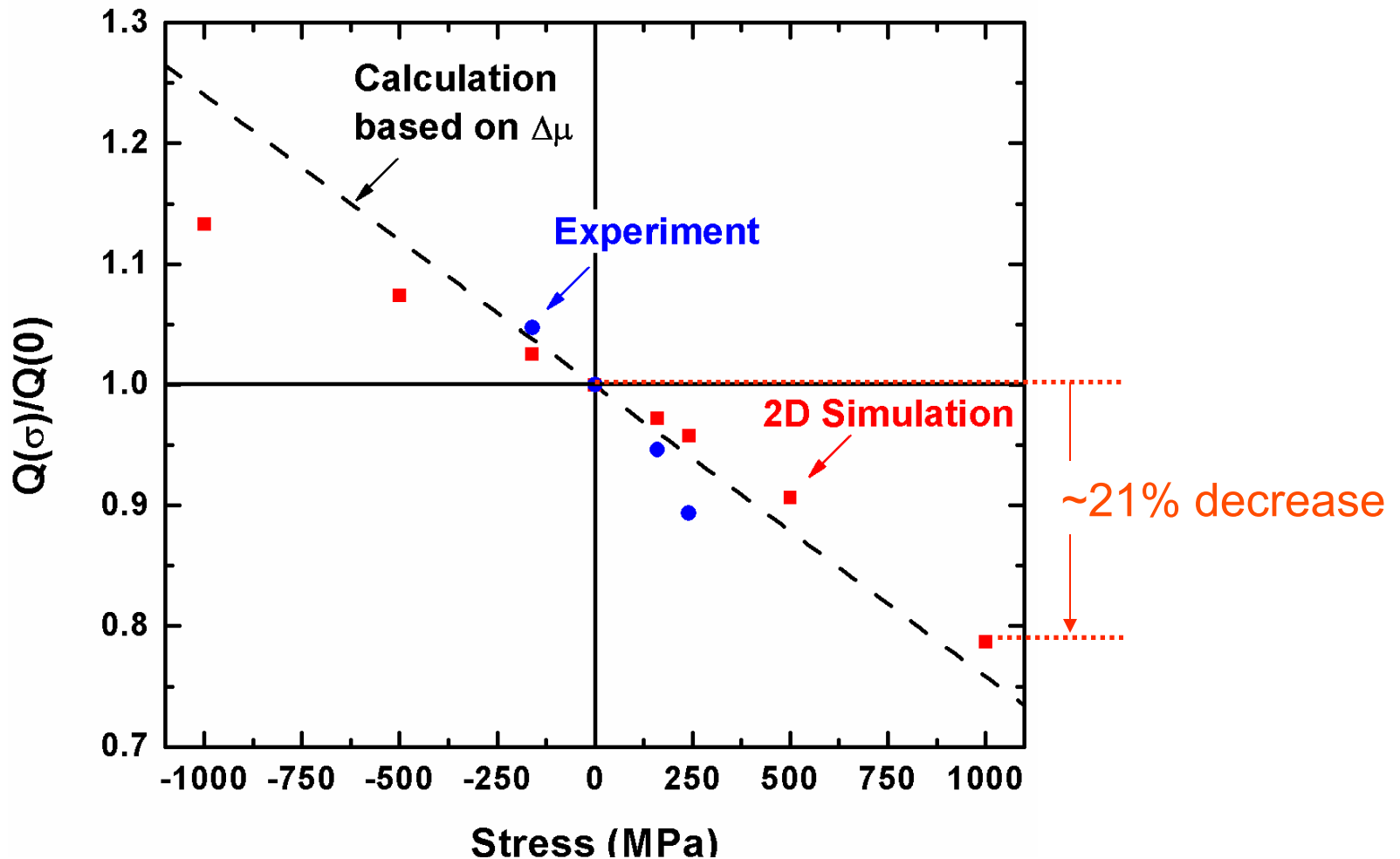
-2D simulation results have good agreement with experimental ones.

I_{max} under Uniaxial Mechanical Stress



- ~ 23% reduction of peak current under 1GPa of uniaxial tensile stress

Charge Collection until 10 ns



- ~ 22% reduction of collected charges under 1GPa of uniaxial tensile stress

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Strain Model (Piezoresistance Effect)



$$\Delta\rho_{zz} / \rho_{zz} = -\Delta\mu_{zz} / \mu_{zz} = \pi_{12}\sigma_{xx} + \pi_{12}\sigma_{yy} + \pi_{11}\sigma_{zz}$$

1) Uniaxial stress along the $\langle 110 \rangle$ direction

$$\Delta\rho_{zz} / \rho_{zz} = -\Delta\mu_{zz} / \mu_{zz} = \pi_{12}\sigma_{xx} / 2 + \pi_{12}\sigma_{yy} / 2 = \pi_{12}\sigma$$

$$\pi_{11} = -102.2 \times 10^{-5} / MPa$$

$$\pi_{12} = +53.4 \times 10^{-5} / MPa$$

2) Uniaxial stress along the $\langle 100 \rangle$ direction

$$\Delta\rho_{zz} / \rho_{zz} = -\Delta\mu_{zz} / \mu_{zz} = \pi_{12}\sigma_{xx} = \pi_{12}\sigma$$

3) Biaxial stress

$$\Delta\rho_{zz} / \rho_{zz} = -\Delta\mu_{zz} / \mu_{zz} = \pi_{12}\sigma_{xx} + \pi_{12}\sigma_{yy} = \pi_{12}(2\sigma)$$

- Mobility along z $\langle 001 \rangle$ -direction change is dominant factor for current transients change.
- Biaxial tensile stress has the potential to decrease $\mu_{n\perp}$ more than uniaxial stresses do.
- It is expected that peak current under biaxial tensile stress is reduced more.

Mobility Enhancement using π - coefficient



$$\begin{aligned}\frac{\Delta\rho_{zz}}{\rho_{zz}(0)} &= \frac{\rho_{zz}(\sigma) - \rho_{zz}(0)}{\rho_{zz}(0)} = \frac{1/\mu_{zz}(\sigma) - 1/\mu_{zz}(0)}{1/\mu_{zz}(0)} \\ &= \frac{\mu_{zz}(0) - \mu_{zz}(\sigma)}{\mu_{zz}(\sigma)} = \pi\sigma\end{aligned}$$

Using power series in $-1 < \pi\sigma < 1$,

$$\mu_{zz}(\sigma) = \frac{1}{1 + \pi\sigma} \mu_{zz}(0)$$

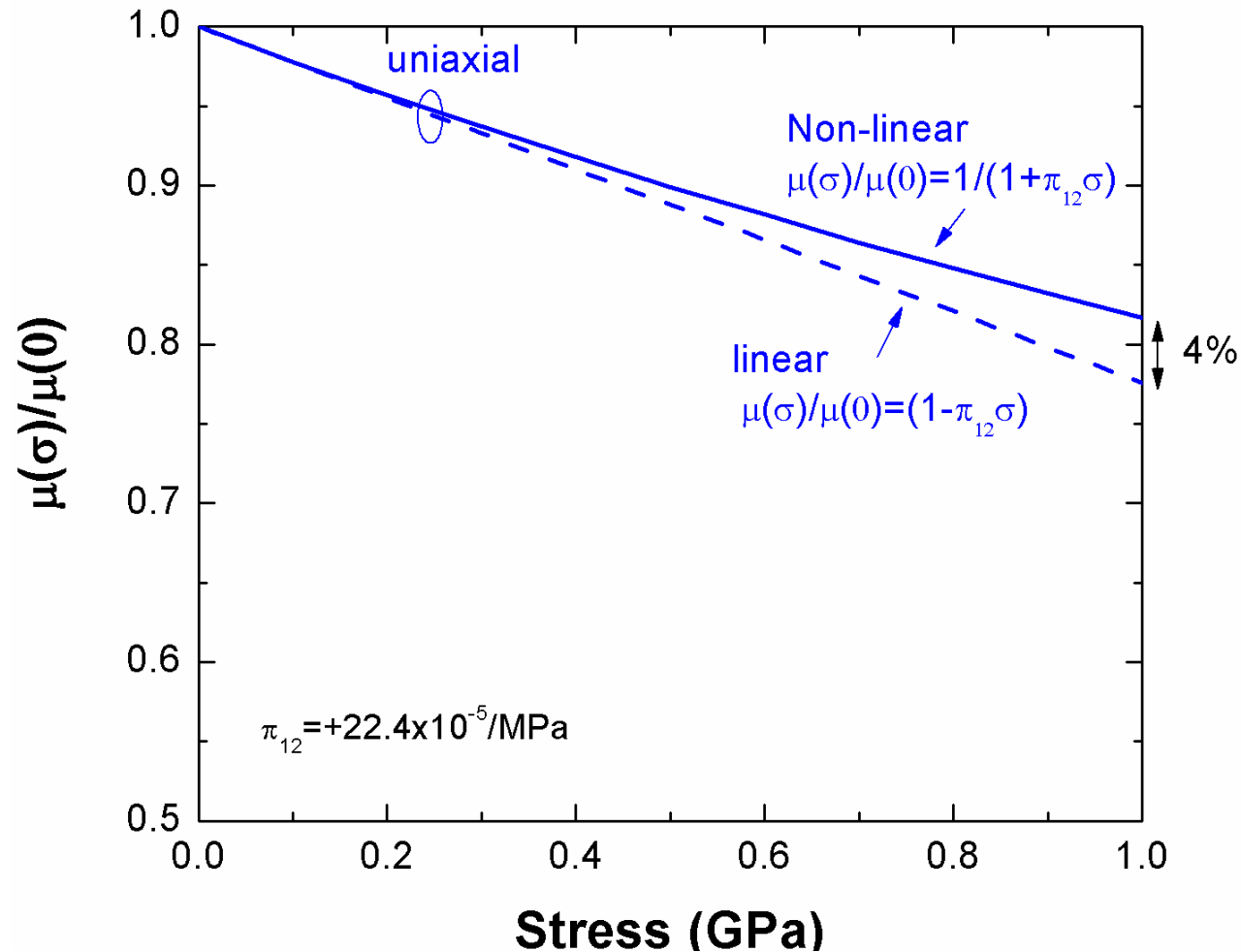
Non-Linear

$$= [1 - \pi\sigma + (\pi\sigma)^2 - (\pi\sigma)^3 + (\pi\sigma)^4 - (\pi\sigma)^5 + \dots] \mu_{zz}(0)$$

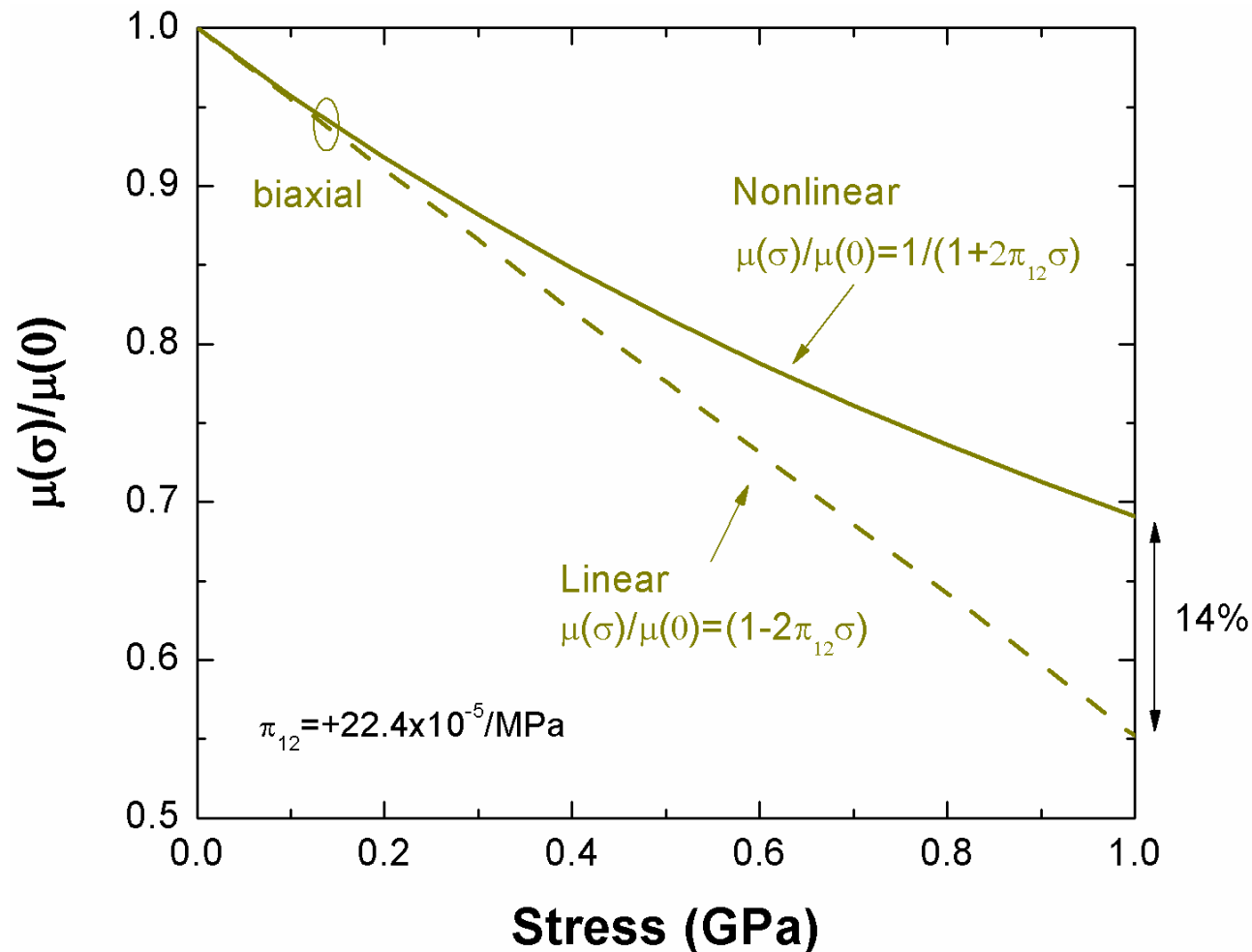
If $\pi\sigma$ is very small, $\mu_{zz}(\sigma) = [1 - \pi\sigma] \times \mu_{zz}(0)$ **Linear**

- We need to consider nonlinear effect in high stress and large piezocoefficient.

Linear vs. Nonlinear in Uniaxial Stress

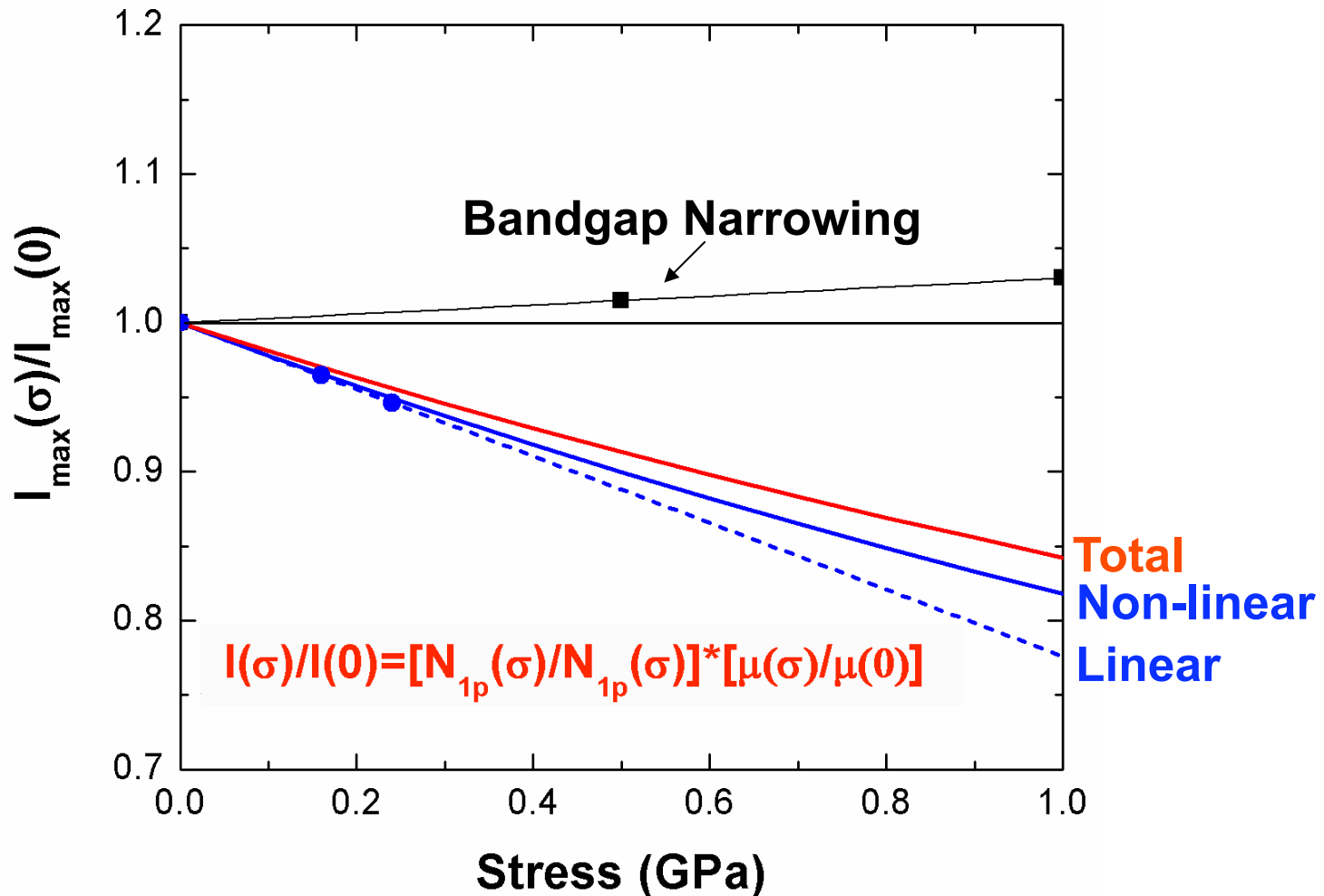


Linear vs. Nonlinear in Biaxial Stress

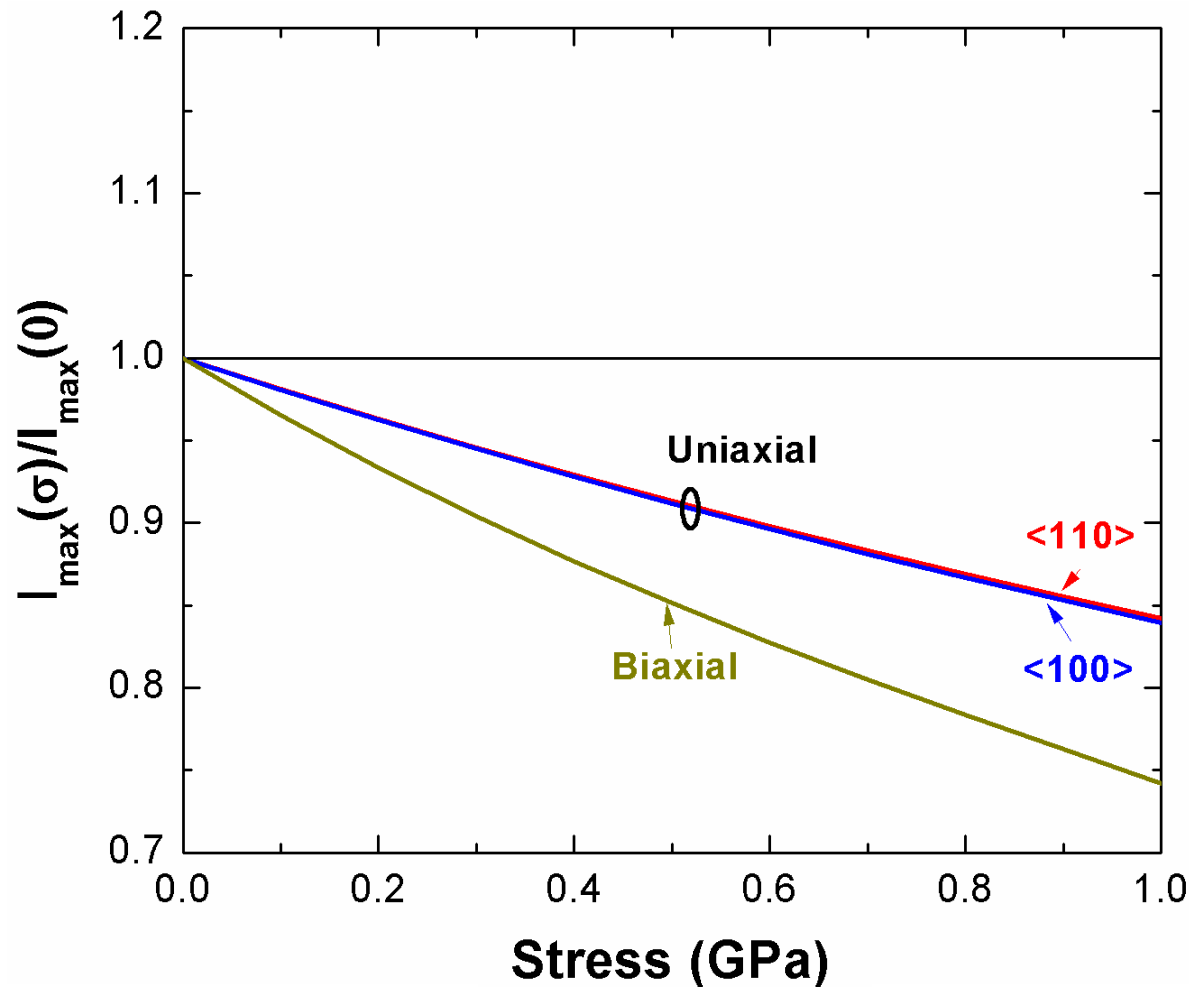


- Biaxial stress shows significant non-linear mobility effect more than uniaxial stress does.

I_{\max} under $\langle 110 \rangle$ Uniaxial Stress



I_{\max} under Uniaxial vs. Biaxial Stress



- Biaxial stress reduces peak current (I_{\max}) more than uniaxial stress does.

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Conclusions/2009 Goals



- Conclusions
 - Experimental results show uniaxial mechanical stress alter current transients due to electron mobility change under stress.
 - 2D simulation results have a good agreement with experiment results in low stress region.
 - Less peak current (~23% decrease at 1 GPa) and charge collection (~21% decrease until 10 ns) in uniaxial tensile stressed diodes are observed.
 - Tensile biaxial stress decreases peak current more than uniaxial stress does.
- Future work
 - Biaxial stress modeling into TCAD simulation
 - Optimum stress type to minimize SET and SEE
 - Hole (P+/n-well) SET modeling

Reference



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Thank you!!

Q & A