

# **Electronic States and Defects in Non-Crystalline and Nano-Crystalline Alternative Transition Metal Delectrics**

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**additional support: DTRA and SRC**

## outline

- (i) high-k dielectrics - 'atomic physics'\*
- (ii) local site symmetries - strong d-state correlation  
Hund's rule - high-spin d-state occupancy
- (iii) relative energies of valence and conduction states  
scale with neutral transition metal atomic energies
- (vi) O-atom vacancy - traps - many electron theory  
**limiting factor - device performance and reliability**  
including radiation hardness

# introduction to many electron theory

## ground and first excited states of the H-atom

Schrodinger equation - neglects electron spin (relativistic effect)

**energy levels**  
from Schrodinger equation agree with Bohr model, and experiment to within *a very small fraction of an eV*

**2s and 2p, etc. at same energy**

$h\nu$  (1s to 2p)  $\sim 13.6 [1-1/4]$   
**10.2 eV or 121.6 nm (UV)**

at high resolution\*  
**i) 2p level split: spin-orbit effect**  
*many electron theory predicts*  
**\*WW2 radar electronics**

**2p<sub>3/2</sub> - 2p<sub>1/2</sub>**

$4.5 \times 10^{-5}$  eV-  $\Delta\nu \sim 10.9$  GHz

$\lambda = 2.7$  cm (micro-wave)

**small difference,**  
**but significant with respect to getting it right**

**many electron wave functions**

***ground state  $1s^1 2p^0$  to excited state  $1s^0 2p^1$***

**term symbol:  $^{2S+1}L_J$**

**L - total orbital angular momentum, S - total spin, and  
J - maximum value of L+S - other J values *stepped down* by 1**

**→ ground state**

**$1s^1 2p^0$ : L = 0, S = 1/2, and J = 1/2**

**$^2S_{1/2}$  with J = +1/2, -1/2, or**

**$^2S_{1/2, -1/2}$  (doubly degenerate)**

**→ excited state**

**$1s^0 2p^1$ : L = 1, S = 1/2, and J = 3/2**

**$^2P_{3/2}$  with J values of +3/2 and +1/2**

**with spin-orbit splitting**

**atomic state symmetries 4- to 8-fold coordination**  
 spherically harmonics -  $SO_3$  symmetry- filled core levels  
 lower symmetries for valence shell electrons

$SO_3$	J	to	$O_h$	$T_d$
s	0	→	$A_{1g}$	$A_1$
p	1	→	$T_{1u}$	$T_2$
d	2	→	$E_g + T_{2g}$	$E + T_2$
f	3	→	$A_{1u} + T_{1u} + T_{2u}$	$A_2 + T_1 + T_2$
g	4	→	$A_{1g} + E_g + T_{2g} + T_{1u}$	$A_1 + E_g + T_2 + T_1$

basis set for Si, Ge, III-V's, Si(GeO<sub>2</sub>), Si(Ge)<sub>3</sub>N<sub>4</sub>, Hf high-k's

$A_{1g}$  and  $B_{1g}$  - non-degenerate d-states -  $dz^2$  and  $dx^2-y^2$

$A_{2g}$  - non-degenerate d-states –  $dxz$  (or  $dyz$ ,  $dxy$ )

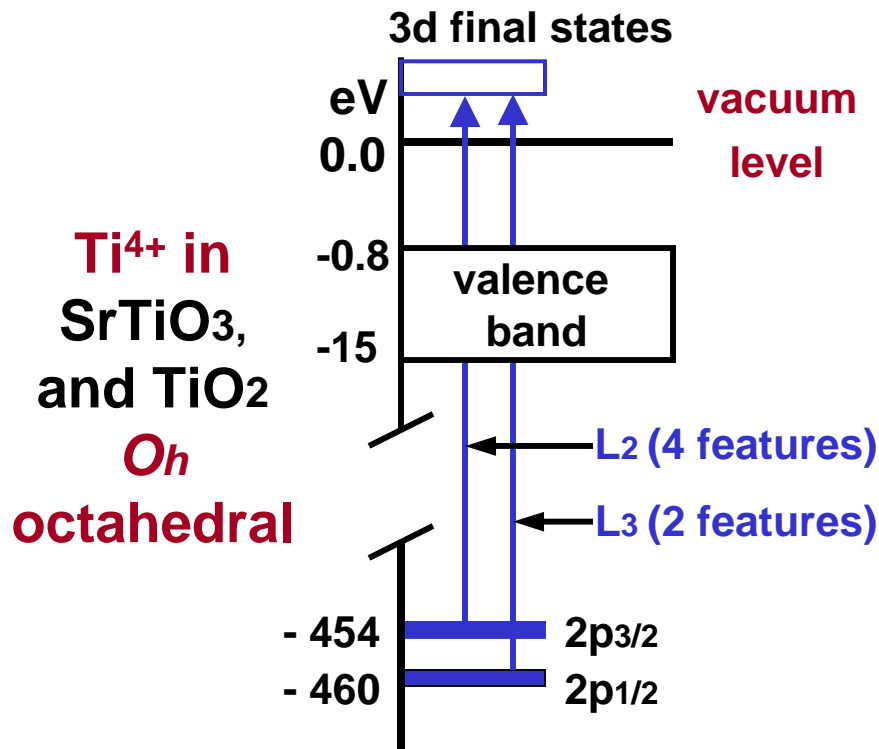
E - doubly degenerate d-states -  $dz^2$ ,  $dx^2-y^2$

$T_1$  - triply degenerate p-states -  $px$ ,  $py$ ,  $pz$

$T_2$  - triply degenerate d-states -  $dxy$ ,  $dxz$ ,  $dyz$

# L<sub>2,3</sub> spectra 2p<sup>6</sup>2d<sup>0</sup> to 2p<sup>5</sup>3d<sup>1</sup> + (2p<sup>6</sup>4s<sup>0</sup> to 2p<sup>5</sup>4d<sup>1</sup>)satellite octahedral Ti<sup>4+</sup>

charge transfer multiplets theory applied to SrTiO<sub>3</sub>\*



ideal octahedral symmetry: Ti<sup>4+</sup> site  
i) number of transitions same for *O<sub>h</sub>* and *SO<sub>3</sub>*, easier to "count in SO<sub>3</sub>

initial state

$$2p^6 2d^0: L = 0, S = 0, J = 0$$

$$2S+1L_J = {}^1S_0 (L = 0)$$

final state

$$2p^5 2d^1$$

$$L = 3, S = -1/2 + 1/2 = 0$$

$$2S+1L_J = {}^1F_{3,2,1} (L = 3)$$

\*F.M.F. de Groot et al, Phys. 42, 928 (1990); Phys. Rev. 42, 5459 (1990).

\*F. de Groot, A. Kotani, *Core level spectroscopy of solids* (Boca Raton, CRC Press, 2008).

## L<sub>2,3</sub> spectra 2p<sup>6</sup>sd<sup>0</sup> to 2p<sup>5</sup>3d<sup>1</sup>

irreducible representation (IR)/ground state; A<sub>1g</sub>  
 A<sub>1g</sub> is a one-dimensional symmetric "d-like" state  
 final states must have "p-like" T<sub>1u</sub> symmetries

$$\langle A_{1g} | T_{1u} \rangle = 0, \text{ equivalent to } \Delta L = \pm 1$$

ground state

SO <sub>3</sub>	J	to	O <sub>h</sub>
s	0	→	A <sub>1g</sub>
p	1	→	T <sub>1u</sub>
d	2	→	E <sub>g</sub> + T <sub>2g</sub>
f	3	→	A <sub>1u</sub> + style="border: 1px solid blue;">T <sub>1u</sub> + T <sub>2u</sub>
g	4	→	A <sub>1g</sub> + E <sub>g</sub> + T <sub>2g</sub> + T <sub>1u</sub>

triply degenerate final  
states  
different *ml* and *j*

3 states

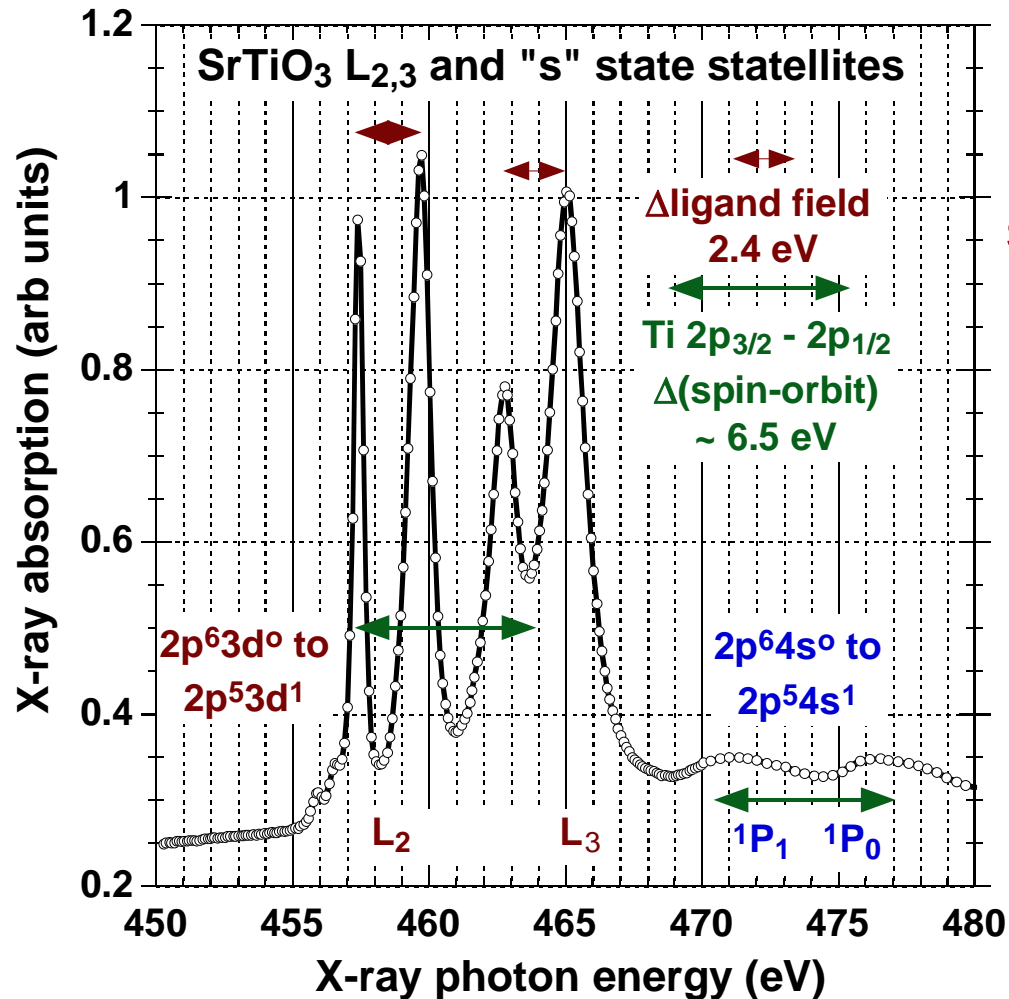
+

3 states

= 6 spectral features

## L<sub>2,3</sub> SrTiO<sub>3</sub>

Ti<sup>4+</sup> – ideal octahedral symmetry, O<sub>h</sub>  
large energy difference between Ti atomic  
3d and 4s states (-11 eV and - 6 eV, respectively)  
expect and observe 6 features



L<sub>2,3</sub> spectra  
2p<sup>6</sup>sd<sup>0</sup> to  
2p<sup>5</sup>3d<sup>1</sup>

4 features in  
L<sub>2</sub> (2p<sub>3/2</sub>)  
2 in L<sub>3</sub> (2p<sub>1/2</sub>)

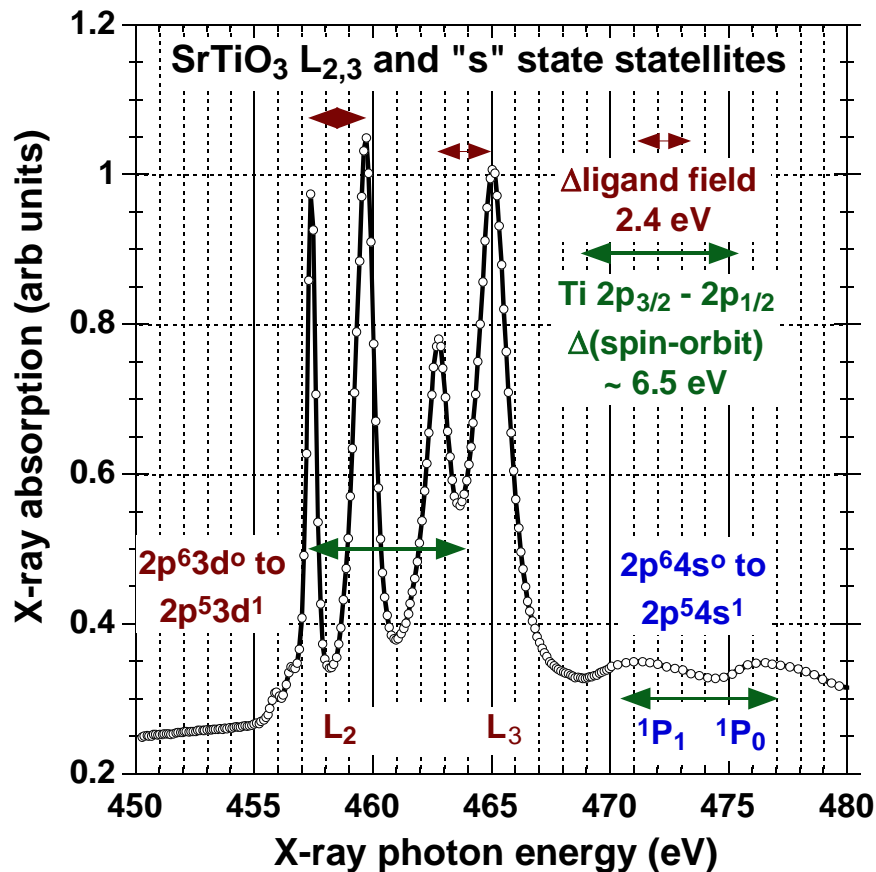
satellite features  
2p<sup>6</sup>4s<sup>0</sup> to 2p<sup>5</sup>4s<sup>1</sup>

X-ray photon  
energies > L<sub>2,3</sub>  
correlates with  
difference of  
5 eV between  
3d and 4s  
Ti atomic states

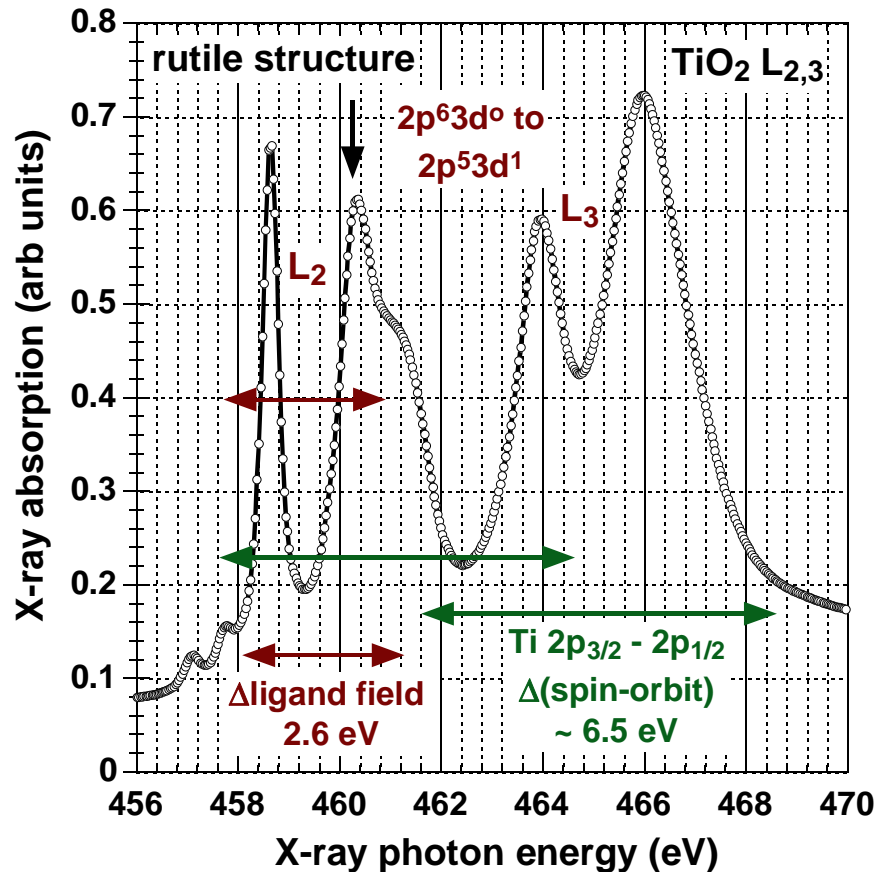


# primitive unit cell distortion in TiO<sub>2</sub> grown on Ge template

elongation of unit cell in z-direction  
removes 2 fold degeneracy of e<sub>g</sub> feature in L<sub>2</sub>

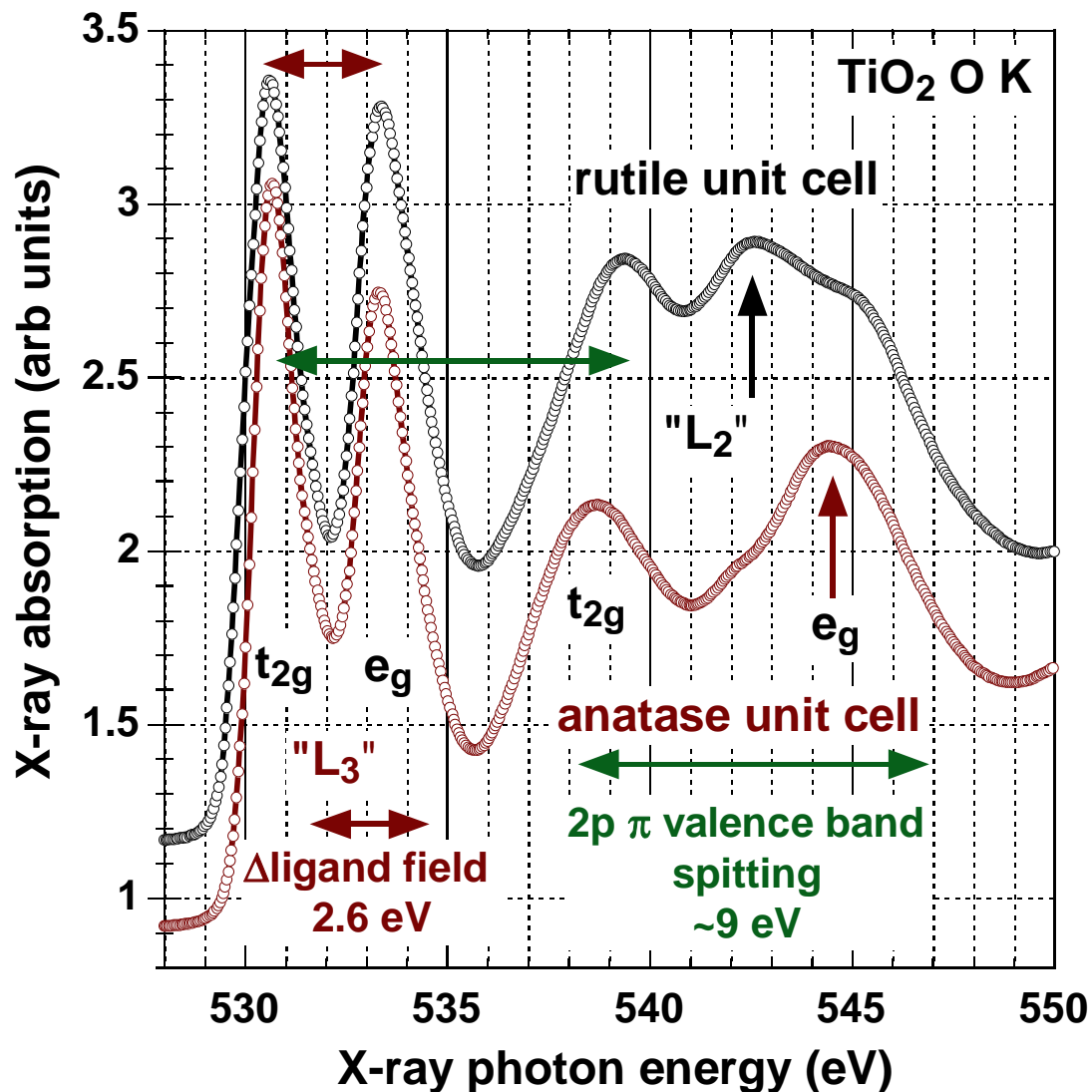


ideal octahedron --  $O_h$  symmetry



tetragonal distortion elongated in z-direction

**differences in L<sub>2,3</sub> between rutile and anatase unit cells**  
**"transported" to O K edge, and then into visible, vacuum ultra-violet (VUV) spectra as well**



(i) corrects previous interpretations of de Groot et al.

**conduction band is not joint DOS from Ti and Partial DOS!!\***

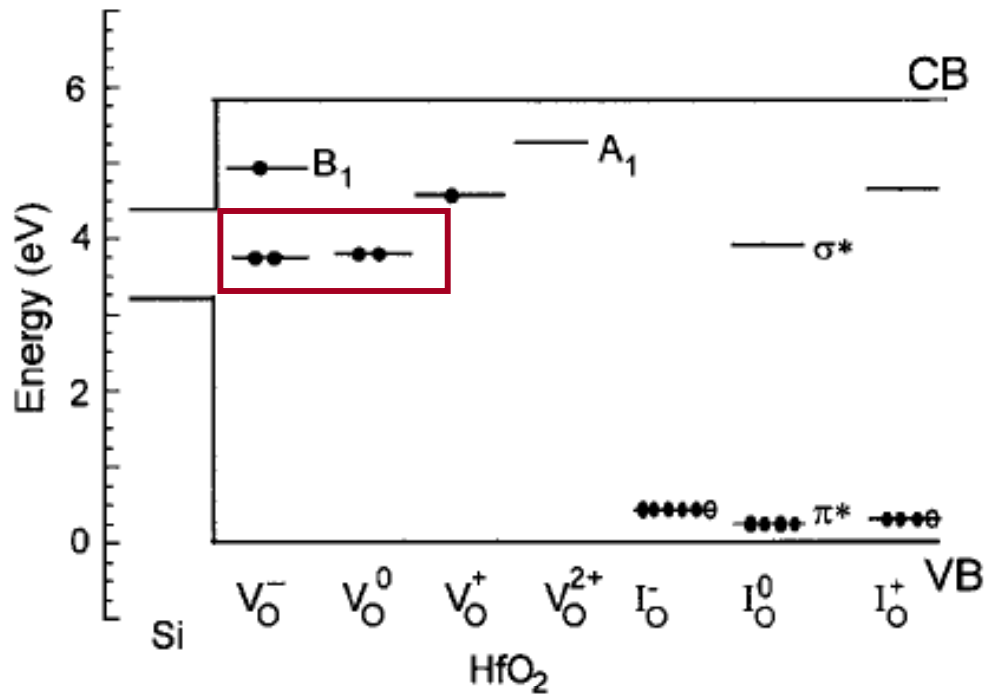
(ii) connection between L<sub>2,3</sub> and OK implications for band edge defects

\* deGroot wrong!!

## **introduction to O-vacancy issues**

**many body theory gives proper interpretation of experiments,  
other theories do not**

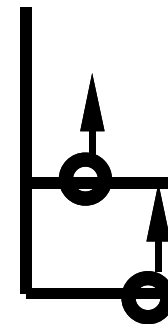
**Robertson, Shlugger, etc...**  
**O-vacancy theory HfO<sub>2</sub>, ZrO<sub>2</sub>**



**many electron theory**  
**ground state**

**<sup>3</sup>F <sup>3</sup>T<sub>1g</sub> J = 1/2, 3/2**

**high spin occupancy**  
**d<sup>2</sup> state**



**schematic**  
**representation of**  
**d<sup>2</sup> occupied states**

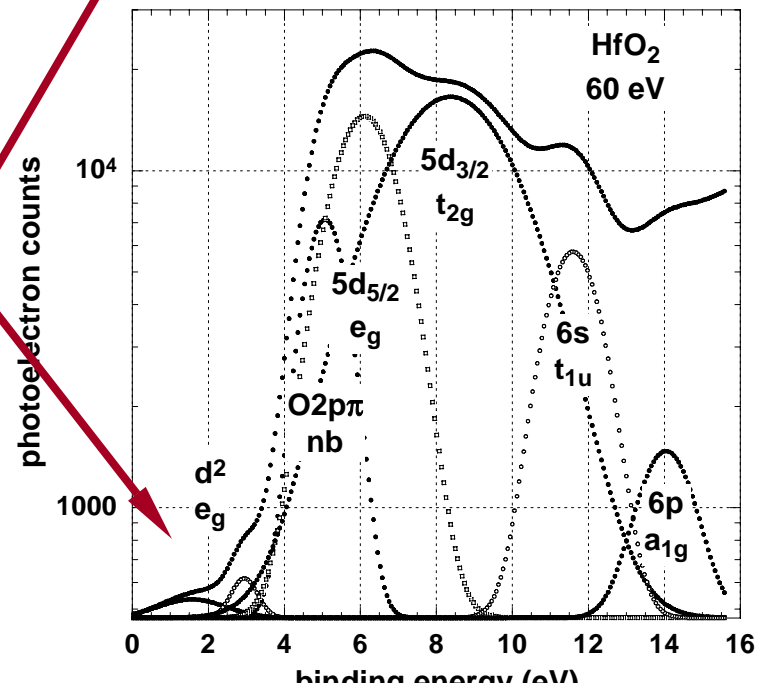
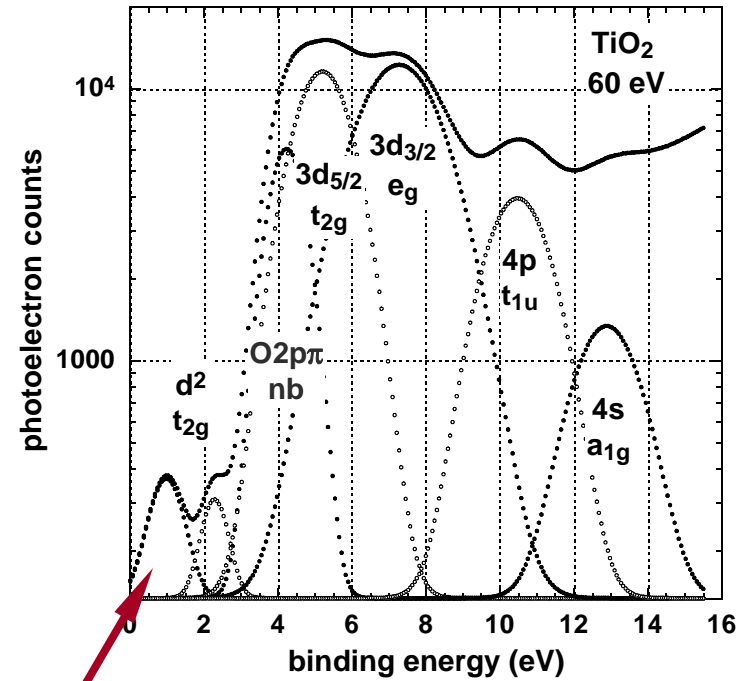
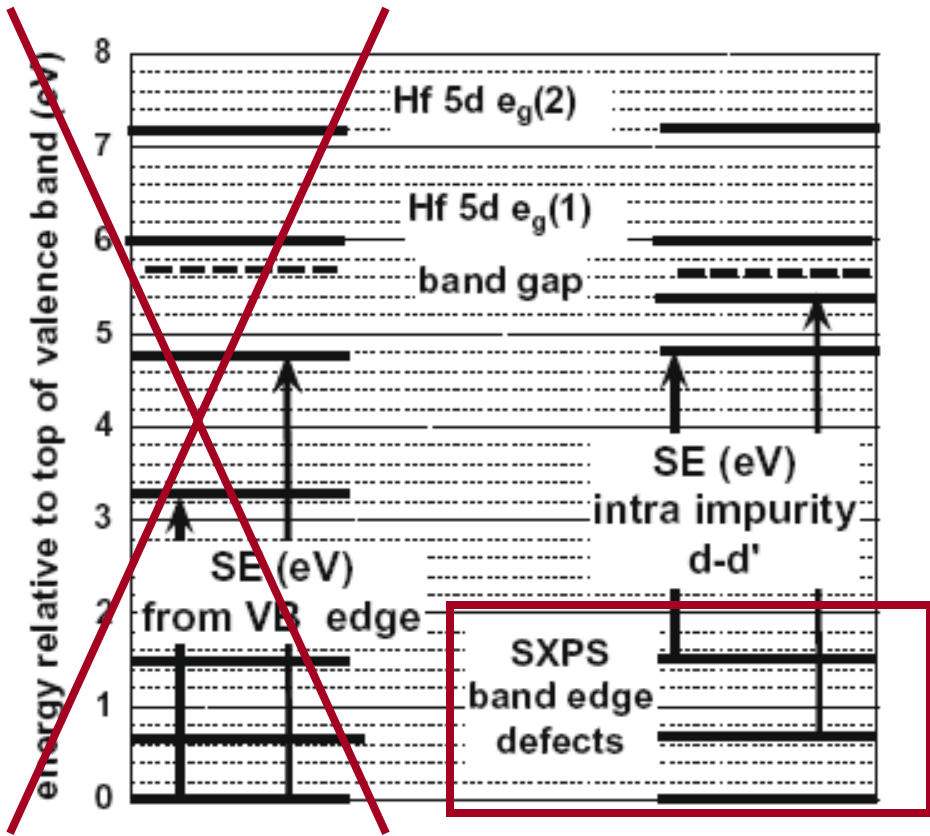
**three serious problems**

➔ **double or low-spin occupancy of Hf(Zr) d-states**  
**R,S ignored Hund's rules, and H Bethe 1929!!**

**encouraged incorrect interpretation of SE,**  
**UCB and ISMT groups**

➔ **ignored NCSU group**  
**valence band spectra**

# UCB and ISMT assigned SE features wrt to top of valence band



**SXPS results confirm R,S had ground state at wrong energy**

## medium range order (SRO) in SiO<sub>2</sub>

MRO associated with d-states - non-degenerate A's, B's, doubly degenerate E's, and triply degenerate p-like T<sub>1</sub>'s and d-like T<sub>2</sub>'s

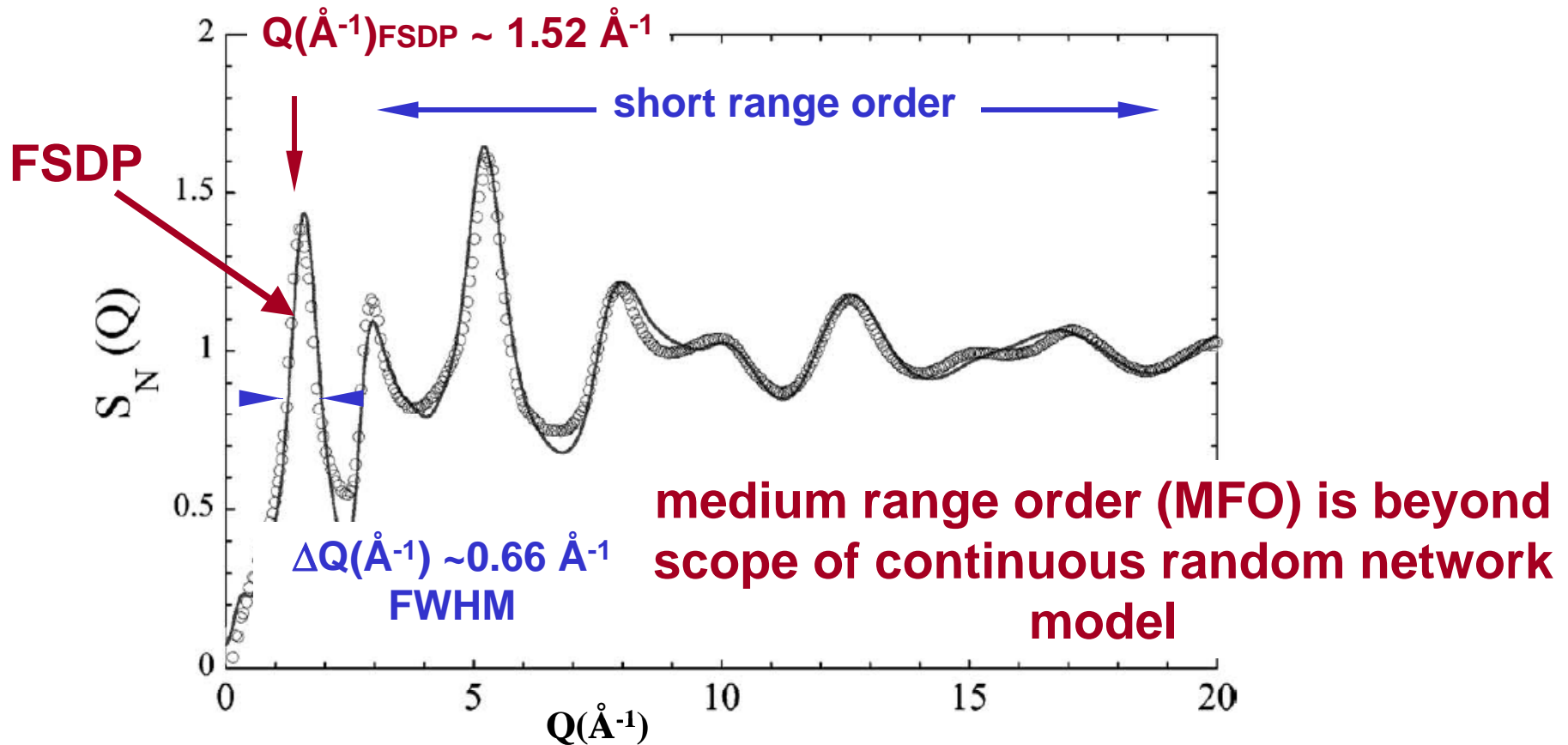
this is what makes SiO<sub>2</sub>, and potentially GeO<sub>2</sub>, and their oxynitrides unique with respect to band edge defect states

prelude to rest of story

major accomplishment is this fundamental intrinsic connection between many electron theory, synchrotron X-ray spectroscopy and defects

- ➔ dielectrics with band edge A<sub>1g</sub> “s-like” states OK, only positive charge after radiation dosing – SiO<sub>2</sub>
- ➔ HfO<sub>2</sub>, etc. with e<sub>g</sub> or t<sub>2g</sub> - O-vacancy/negative ion electron traps – negative charge after radiation dosing

# first sharp diffraction peak\* in structure factor, $S_N$ , for $\text{SiO}_2$



$S_N(Q)$  is Fourier transform of diffraction intensity  
two MRO length scales

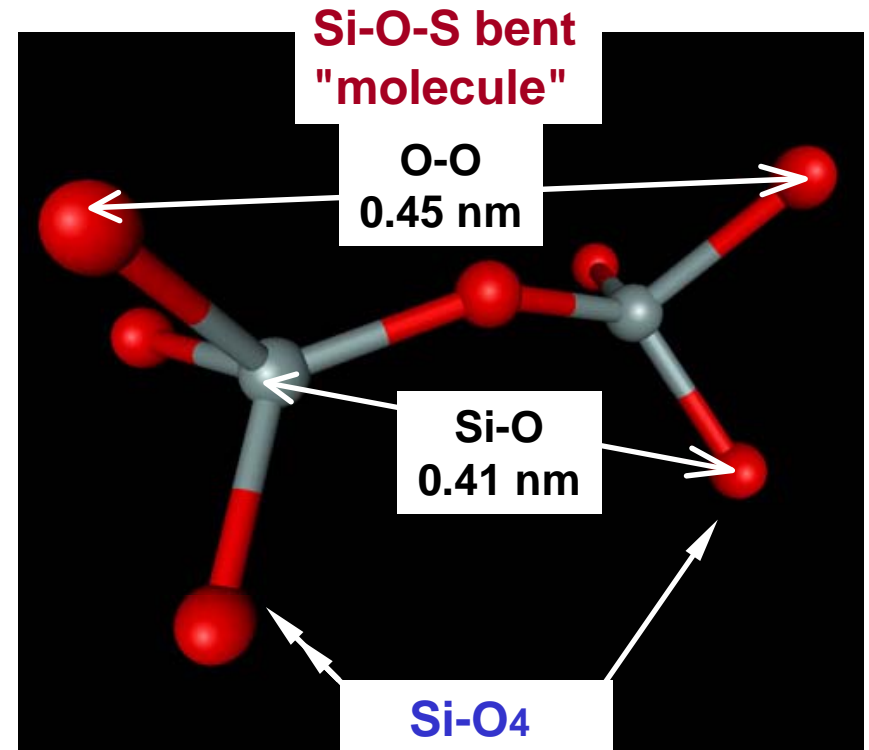
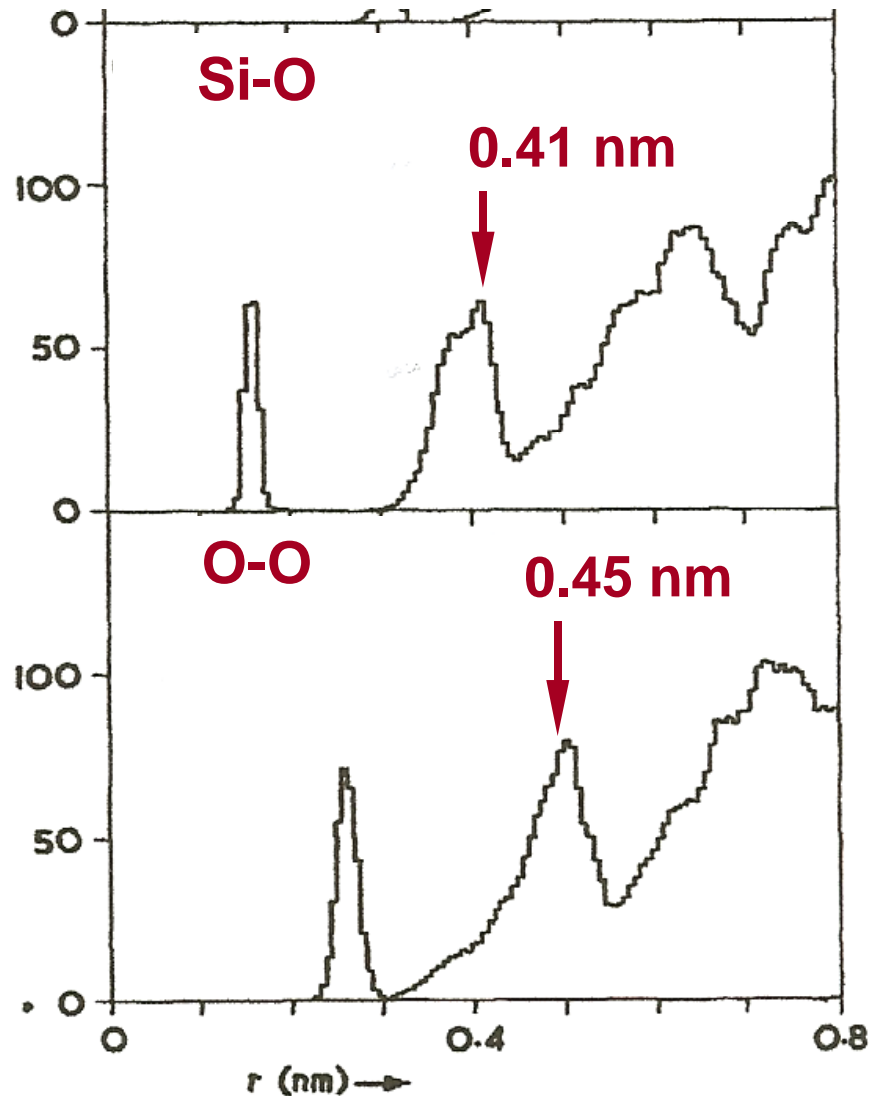
peak position: correlation length,  $\lambda_{\text{corr}} \sim 2\pi / Q(\text{Å}^{-1})$ ,  $\sim 0.4 \text{ nm}$

FWHM: coherence length  $\sim \lambda_{\text{coh}} \sim 2\pi / \Delta Q(\text{Å}^{-1})$ ,  $\sim 0.95 \text{ nm}$

\*J Du & LR Corrales, Phys Rev B 72, 092201 (2005).

**$S_N(Q)^*$  and Bell and Dean\*\* saw it “brushed it off!!”**

**Si-O 3rd and O-O 4th neighbor distances not predicted by CRNs**



$\lambda_{\text{corr}} \sim 2\pi / Q(\text{\AA}^{-1}), \sim 0.4 \text{ nm}$   
intra-six member ring  
 $\lambda_{\text{coh}} \sim 2\pi / \Delta Q(\text{\AA}^{-1}), \sim 0.95 \text{ nm}$   
connected 6-member rings

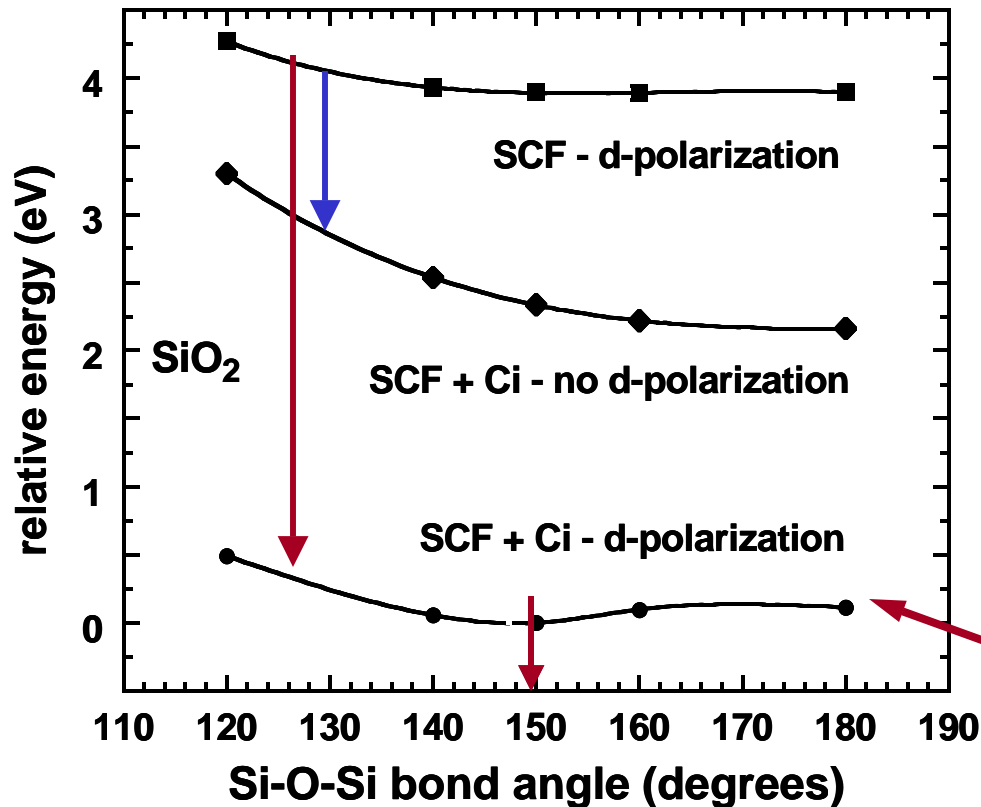
**\*J Du & LR Corrales, Phys Rev B 72, 092201 (2005).**

**\*\*RJ Bell and P Dean, Philos. Mag. 25 (No. 6) 1381-1398 (1972).**



## relative energy vs Si-O-Si bond angle

Si-Si distance fixed - 0.31 nm - O'Keefe and Hyde (1978)



Hartree-Fock (SCF)  
with gaussian  
d-orbitals (polarization)  
on Si

SCF + configuration  
interaction (CI)  
optimization  
(no d-orbitals)

SCF + CI with  
d-orbitals

SiO<sub>2</sub> - 19±2° i) N-L\* 17±2° ii) M-W\*\* >30°

GeO<sub>2</sub> - 11.5±1.5° i) N-L\* 15±2°

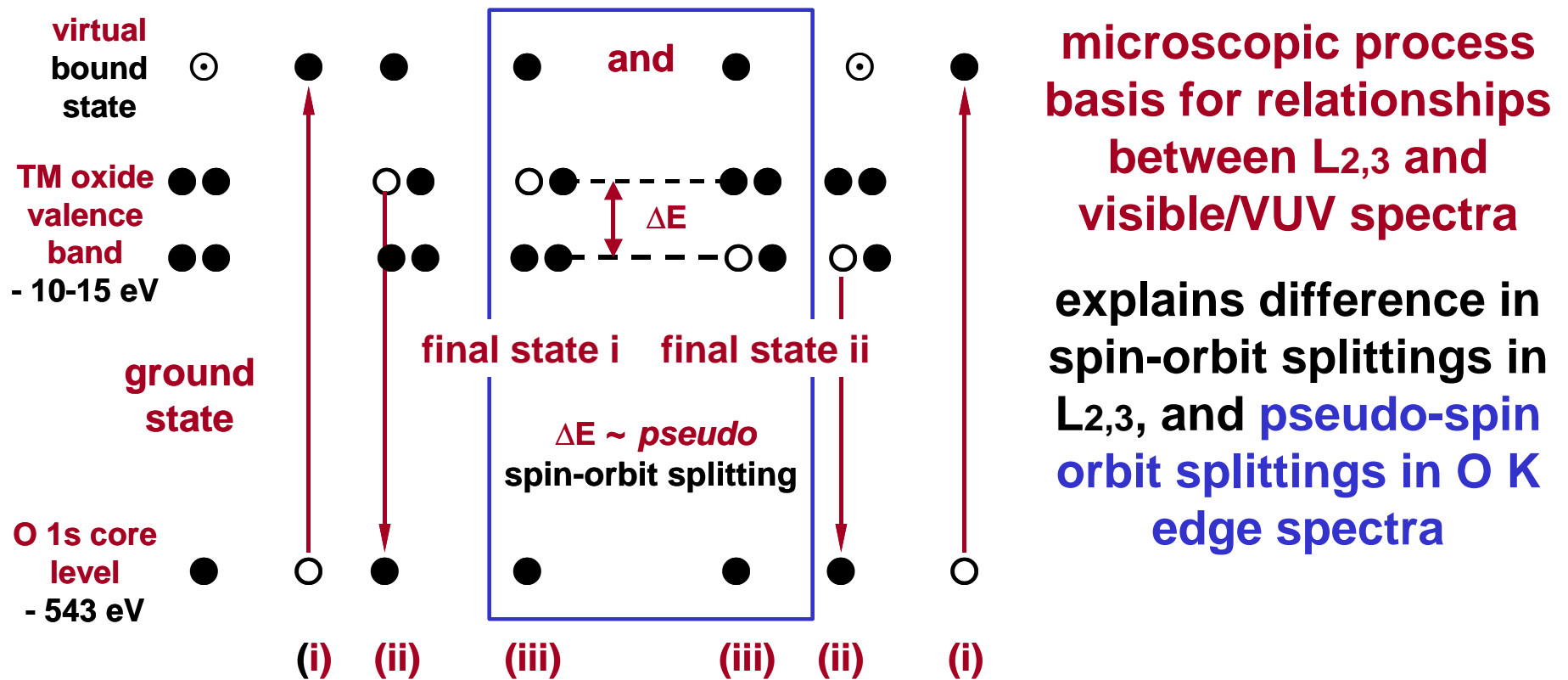
i) N-L\* - J. Neufeind and K.-D. Liss, *Bur Bunsen Phys Chem* 100, 1341 (1996).

ii) M-W\*\* - R.L. Mozzi and B.E. Warren, *J. Appl. Cryst.* 2, 164 (1969).

**O K edge excitations - coherent process for SiO<sub>2</sub> OK edge**

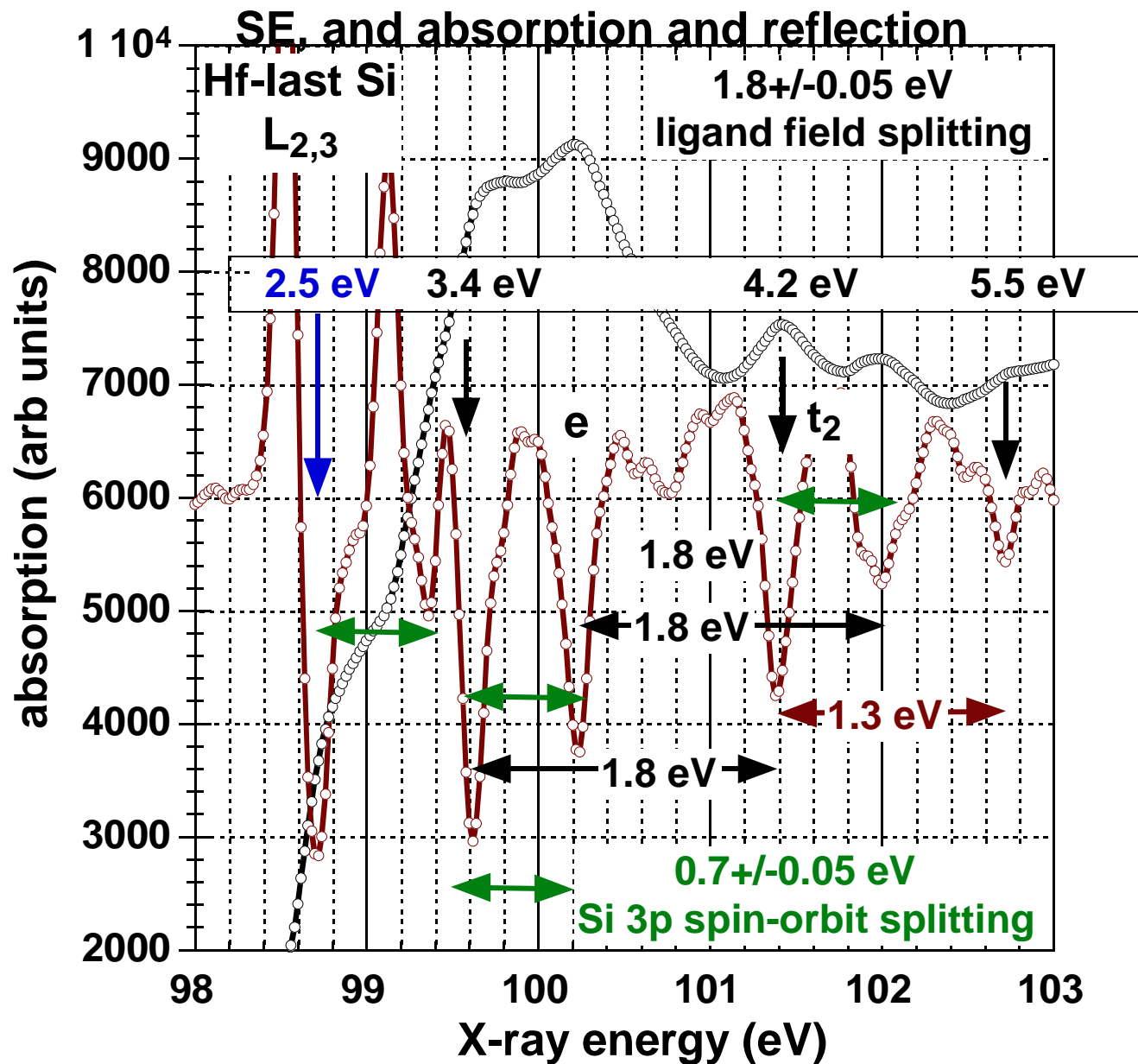
**(i) x-ray excitation to to virtual bound state and**

**(ii) 1s core hole filling form O 2p  $\pi/\sigma$  valence band states**

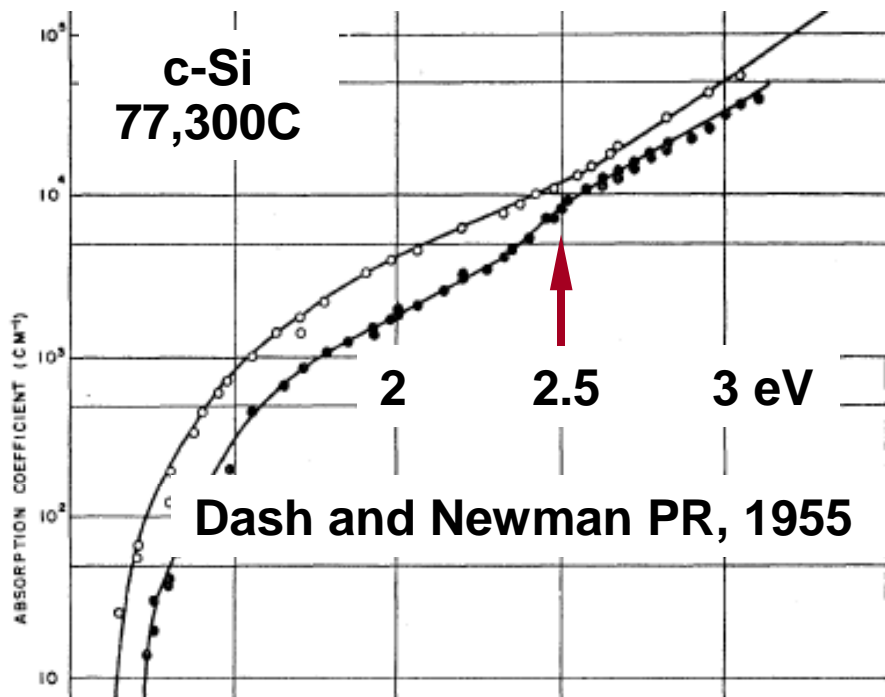


**for Si L<sub>2,3</sub> crystal field splitting of 1.8 eV is greater than S-O of 0.7 eV, whereas for ZrO<sub>2</sub> - spin-orbit, and pseudo spin orbit (13 eV and 10 eV) are greater than crystal field, ~3.5 eV**

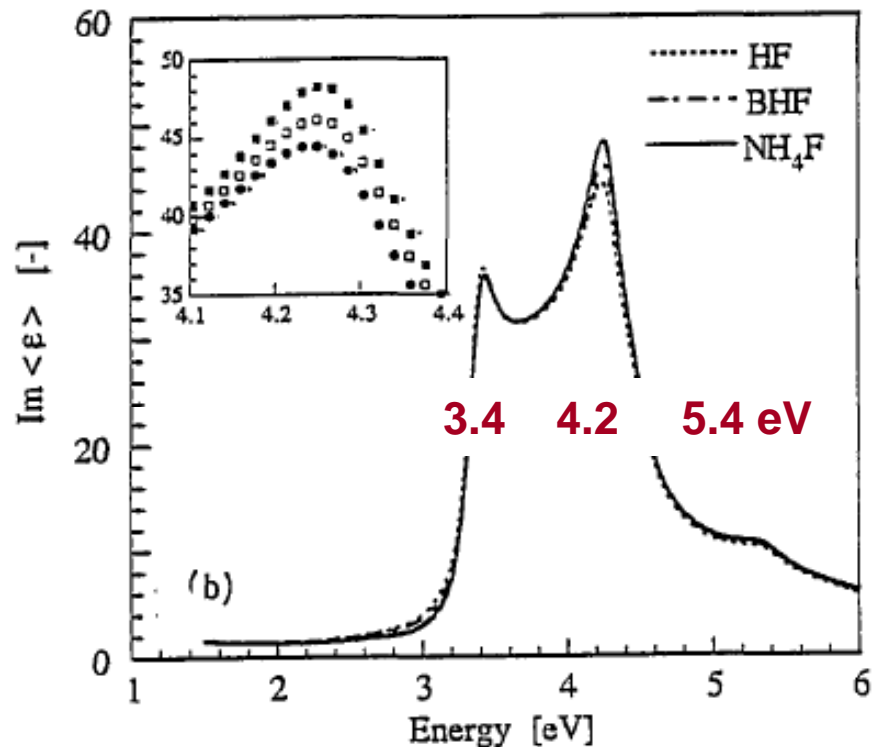
**relative energies of features in 2nd derivative of Si L<sub>2,3</sub>**  
when "transported" to visible/VUV energy scale, correspond to features in



**relative energies of features in 2nd derivative of Si L<sub>2,3</sub>**  
 when "transported" to visible/VUV energy scale, correspond to  
 features in SE, and absorption and reflection



**2.5 eV is direct gap**  
**1.1 eV is indirect gap**  
**phonon-assisted for symmetry**

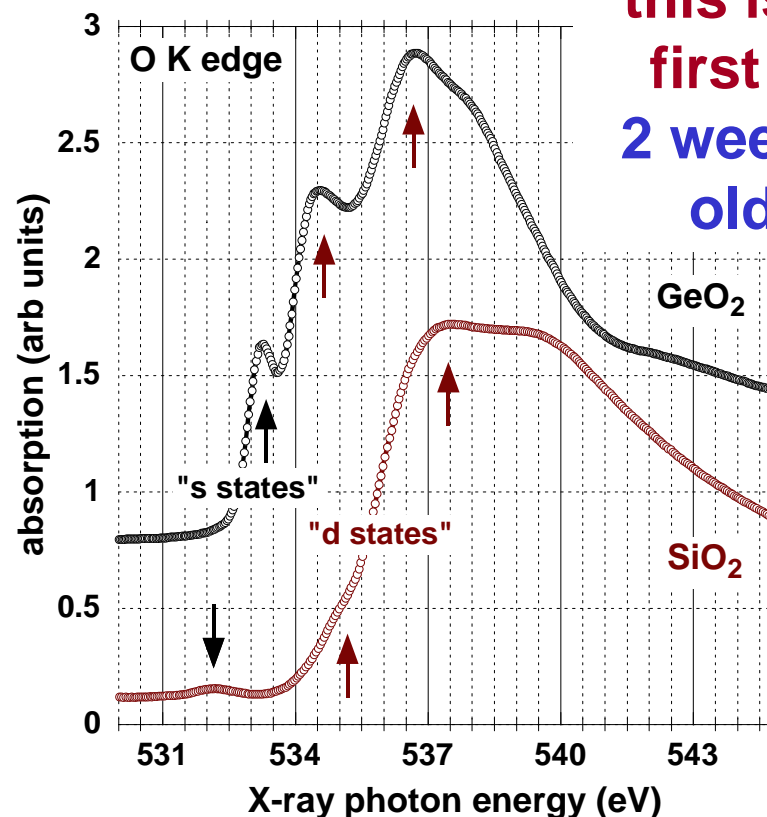
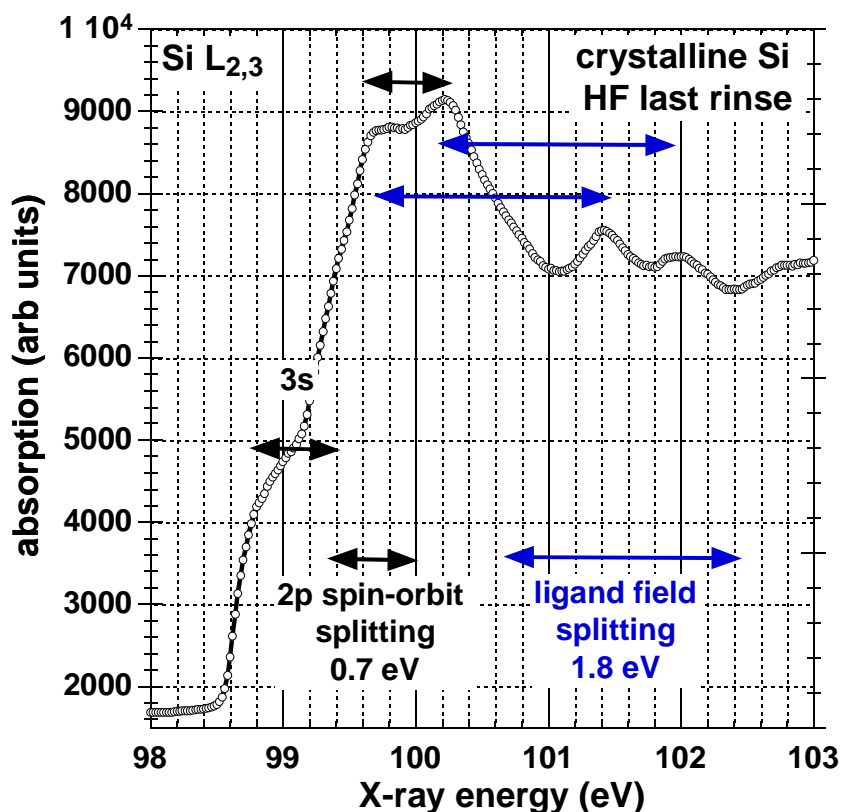


**stronger transitions – d-states!!**  
**in reflection spectroscopic**  
**ellipsometry**

Yasuda and Aspnes JOSB, 1994

# relevance to electrical performance and reliability, including rad-hardness of SiO<sub>2</sub>

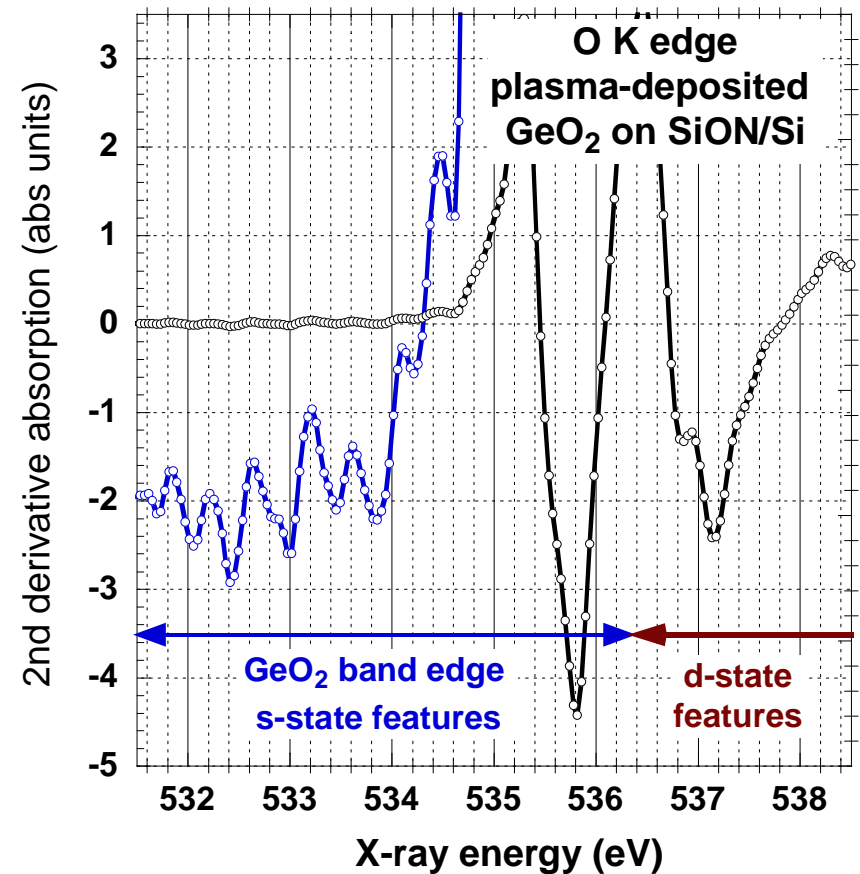
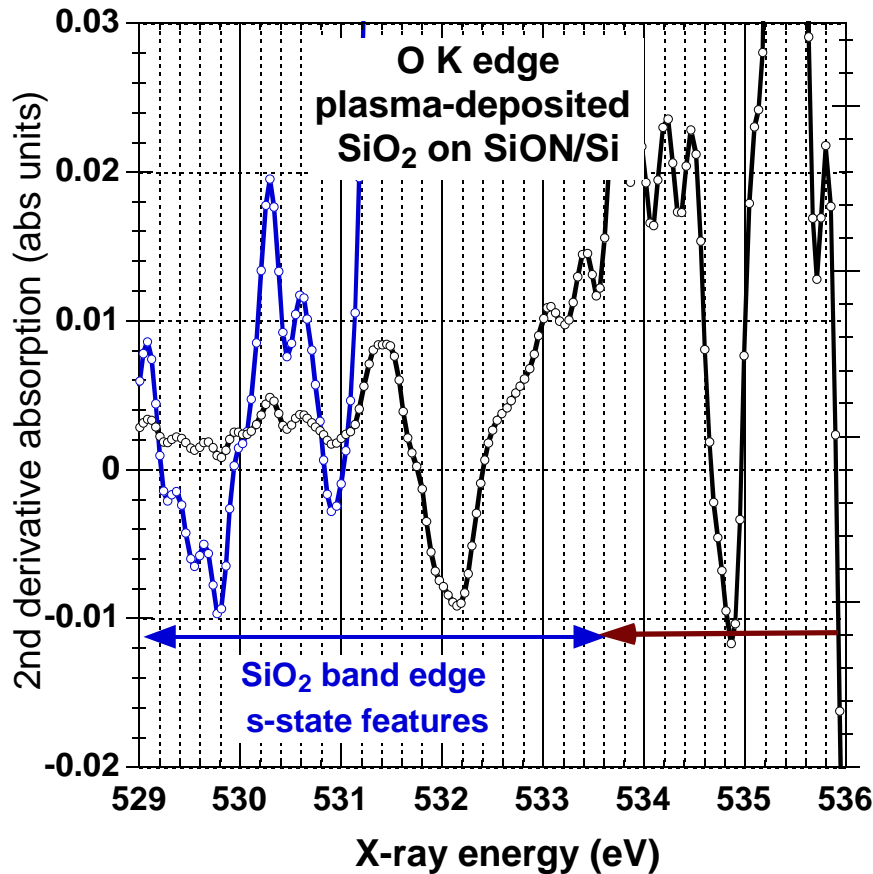
Si "s-like" one-dimension A<sub>1g</sub> features are at  
conduction edges of crystalline-Si and non-crystalline SiO<sub>2</sub>  
and remote plasma deposited GeO<sub>2</sub> (a first!!)



this is a  
first !!  
2 weeks  
old

“s-states” SiO<sub>2</sub>/GeO<sub>2</sub> more stress resistant than Hf-or Lu, La based  
high-k dielectrics – in which band edge states have d-character

relative energies of "s-state" band edge and higher energy "d-state" features are approximately the same in SiO<sub>2</sub> and GeO<sub>2</sub>



**O-atom vacancy defects are in middle of gap!!**

**this makes SiO<sub>2</sub>/GeO<sub>2</sub> special**

**no band edge traps, or negative trapping after X-ray,  $\gamma$ -ray stress**

## **application to experimental results**

### **i) defect suppression in**

**thin HfO<sub>2</sub> - compare 2 nm and 4 nm thick films**

**relationship to trap-assisted tunneling**

### **ii) defects in HfSiON ternaries**

**compare with C-V characteristics**

## reminder: soft X-ray photoelectron spectroscopy valence band spectrum for nano-grain HfO<sub>2</sub>

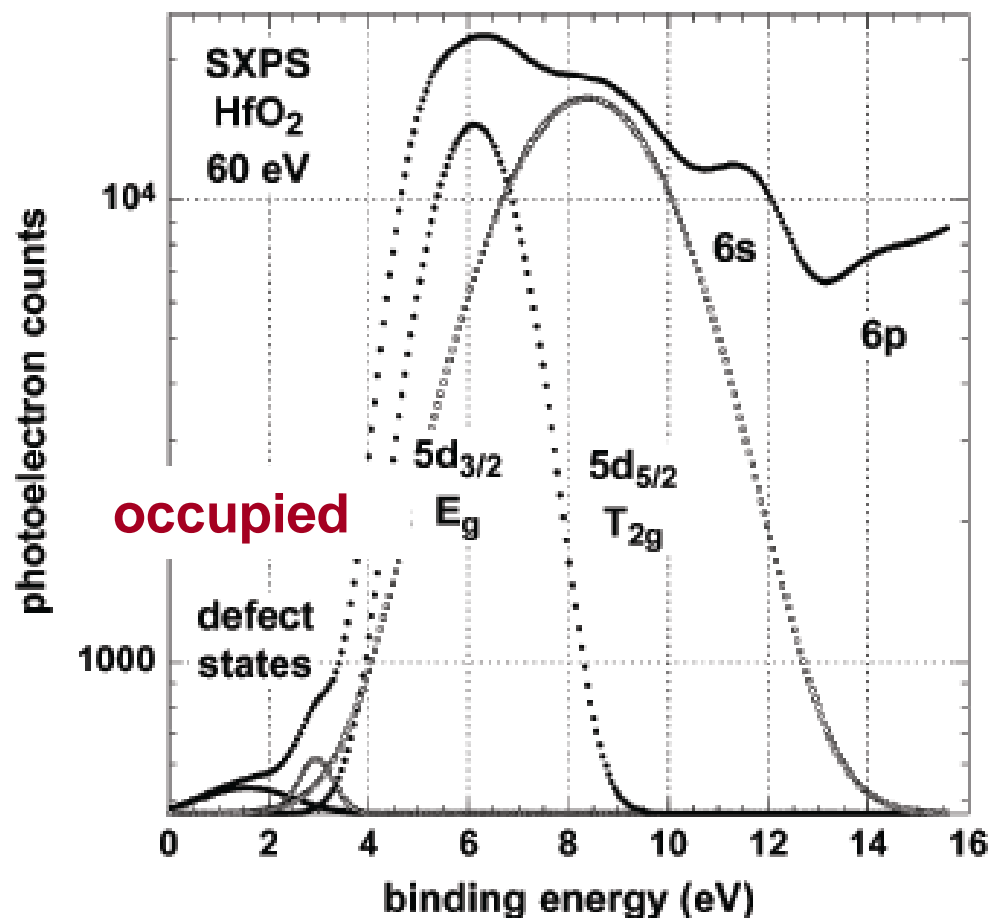


Fig. 5. SXPS valence band spectrum for nano-crystalline HfO<sub>2</sub>: 5d-state and defect features are identified.

O-vacancy electron states – 2 electrons/removed neutral O-atom

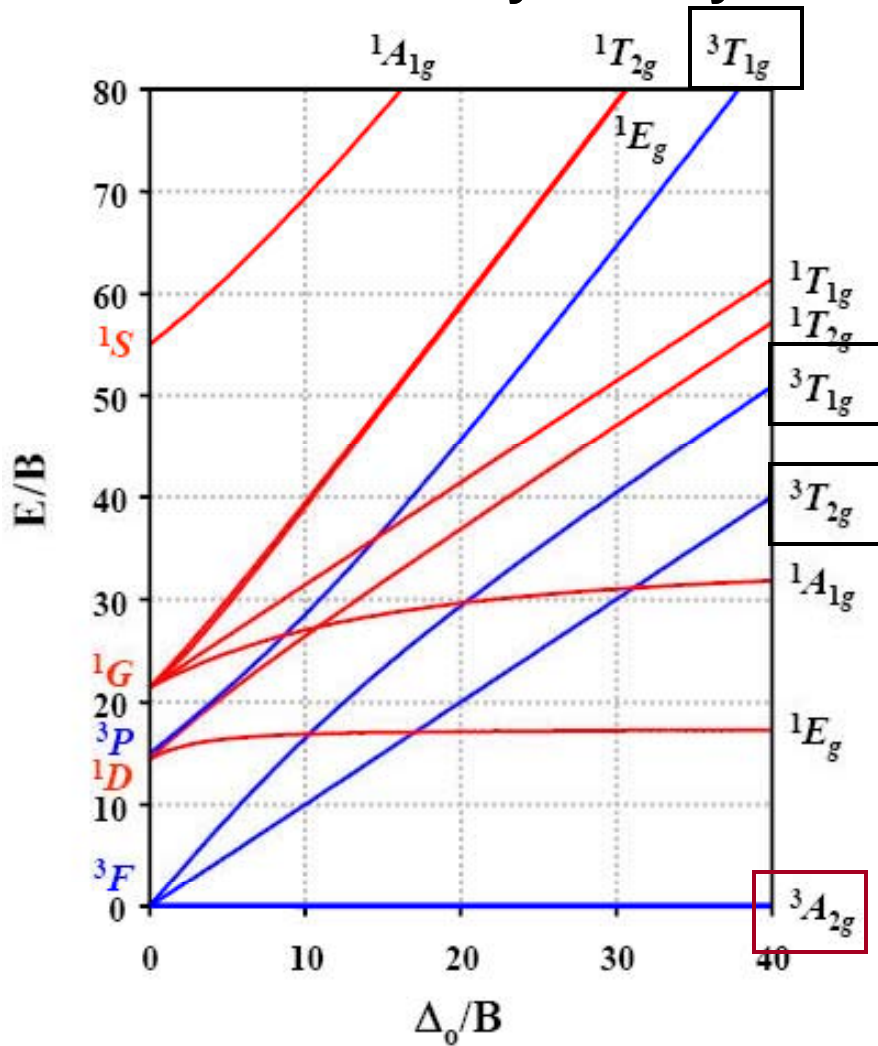
lowest energy states is high spin d<sup>2</sup> that can be on any one of four Hf atoms that border vacancy

excited states (d to d') are near conduction band are electron traps



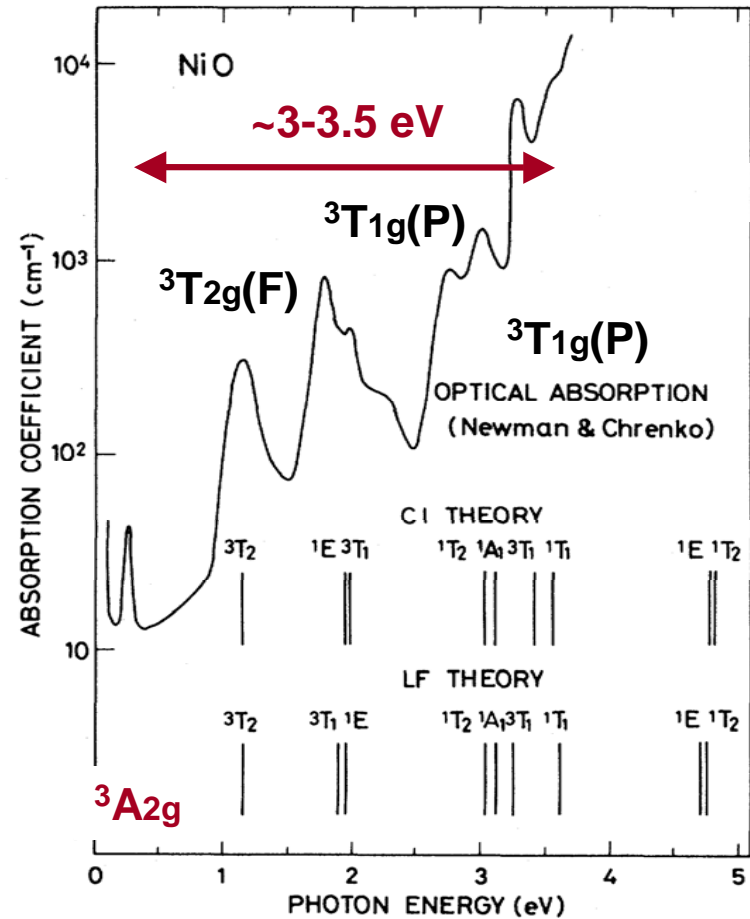
# Tanabe-Sugano diagram

## $d^2$ state occupancy and $O_h$ cubic symmetry



ground triplet 's' state  $3A_{2g}$  to 3 triplet excited 'p'-states,  $3T_{2g}$ ,  $3T_{1g}$  and  $3T_{1g}$

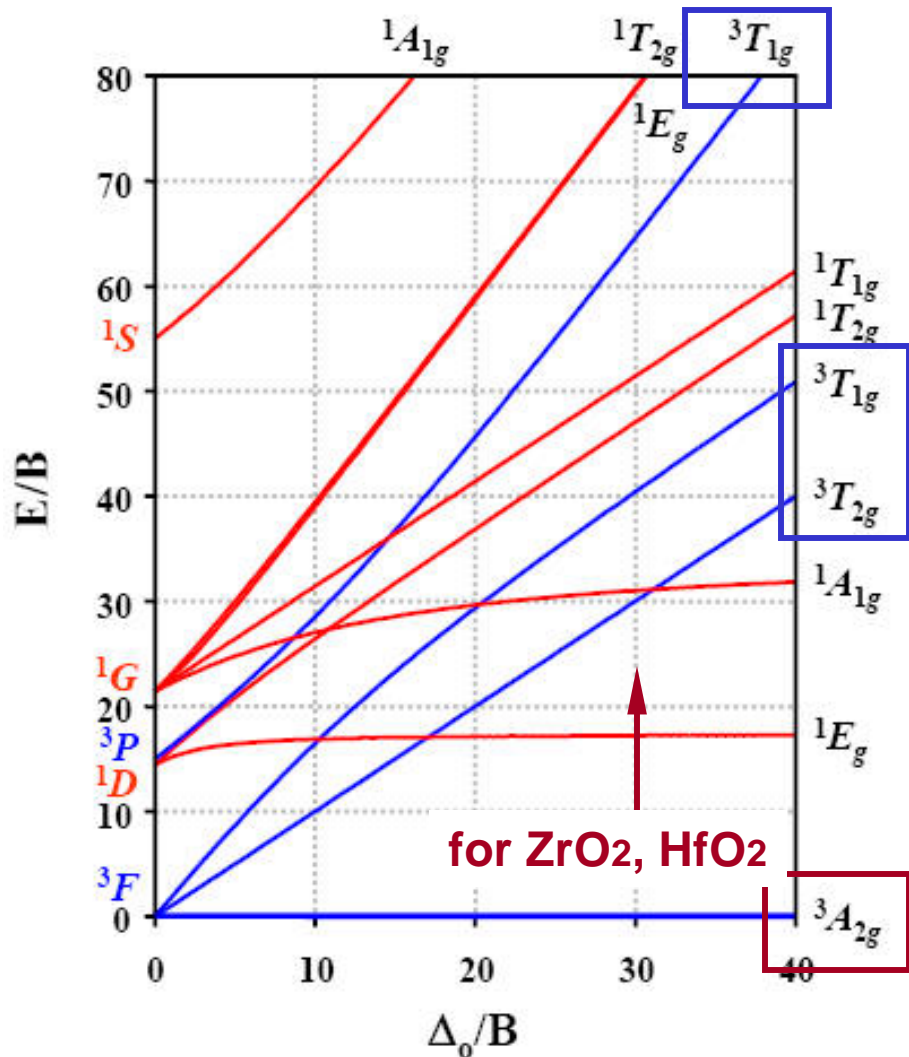
# NaCl structured NiO



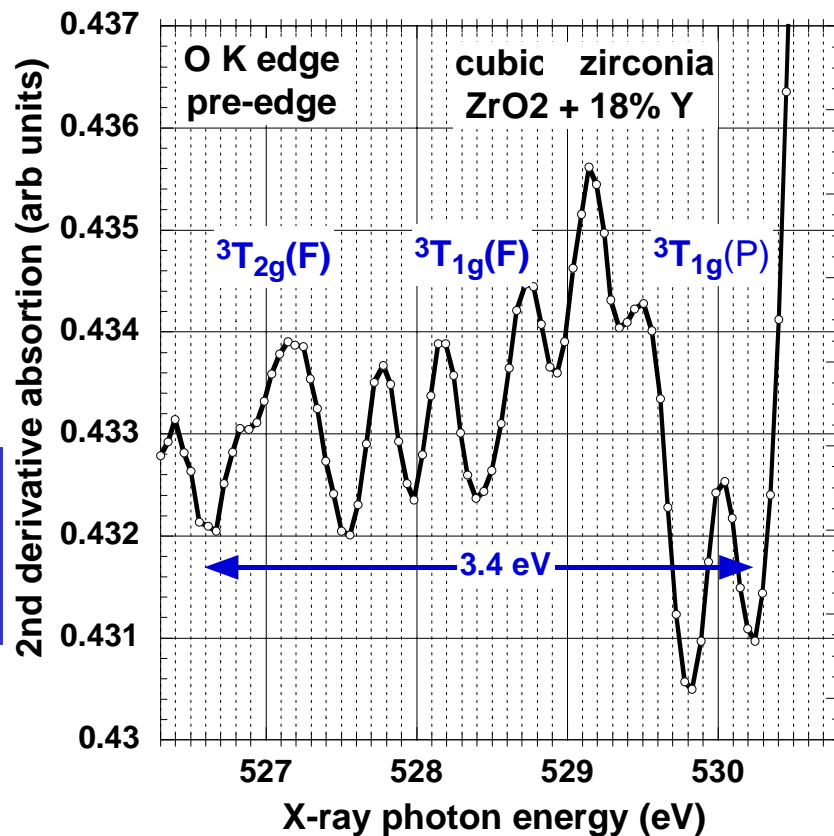
Tanabe-Sugano diagrams also apply to d-d' transitions for O-atom vacancy defects, and negative ion states

A Fujimori, et al., PR B30, 967 (1984).

same Tanabe-Sugano diagram and energy scale applies to  $d^2$  O-atom vacancies in Y, Or, Hf oxides, and NiO in NaCl crystals

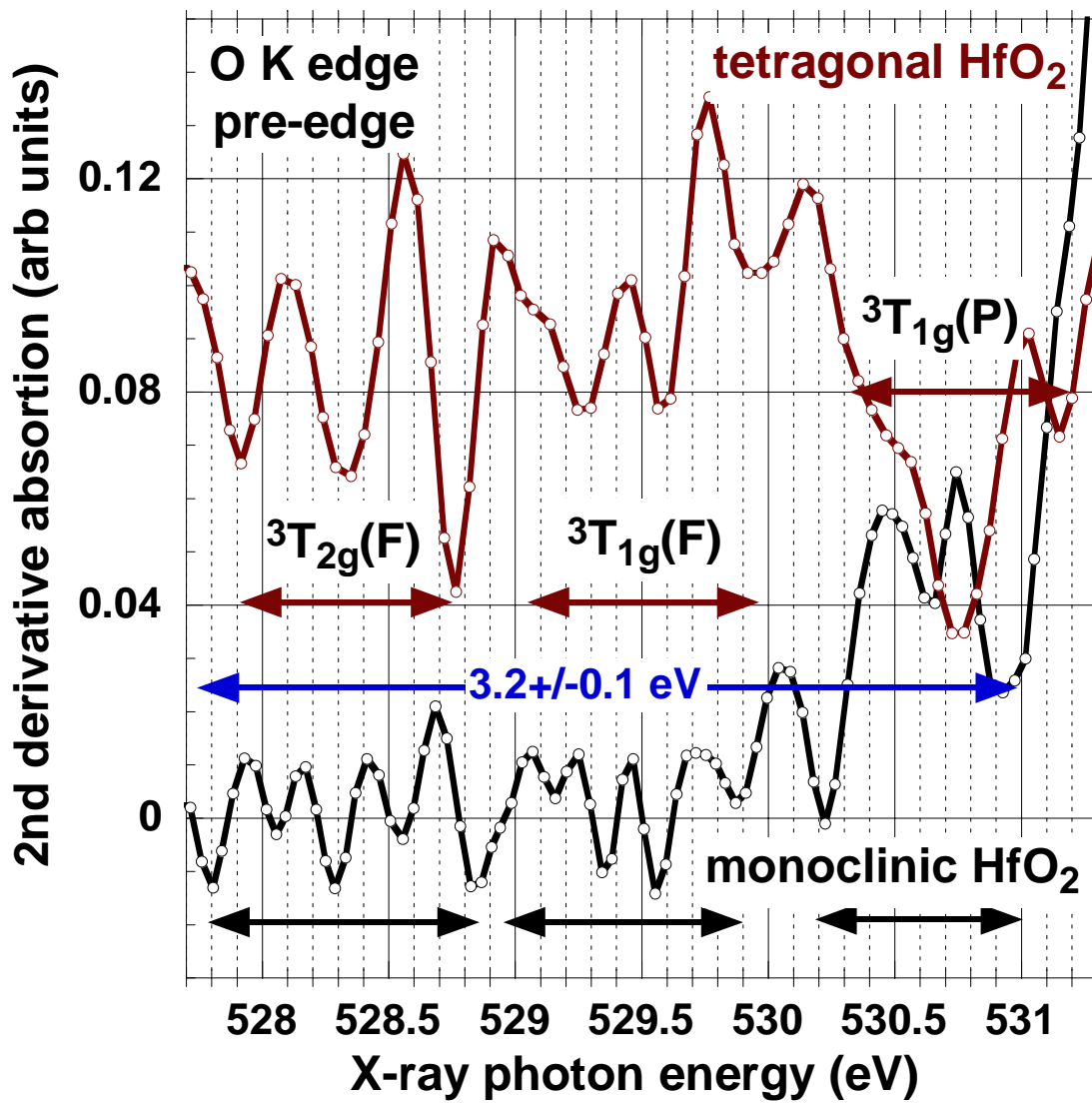


$E/B \sim 60-70 \sim 3-3.5 V$



triplet character of states  
 3 three states with small splitting,  $\sim 0.5$  eV,  
 different values of J, also  
 $\sim$ same spectral width

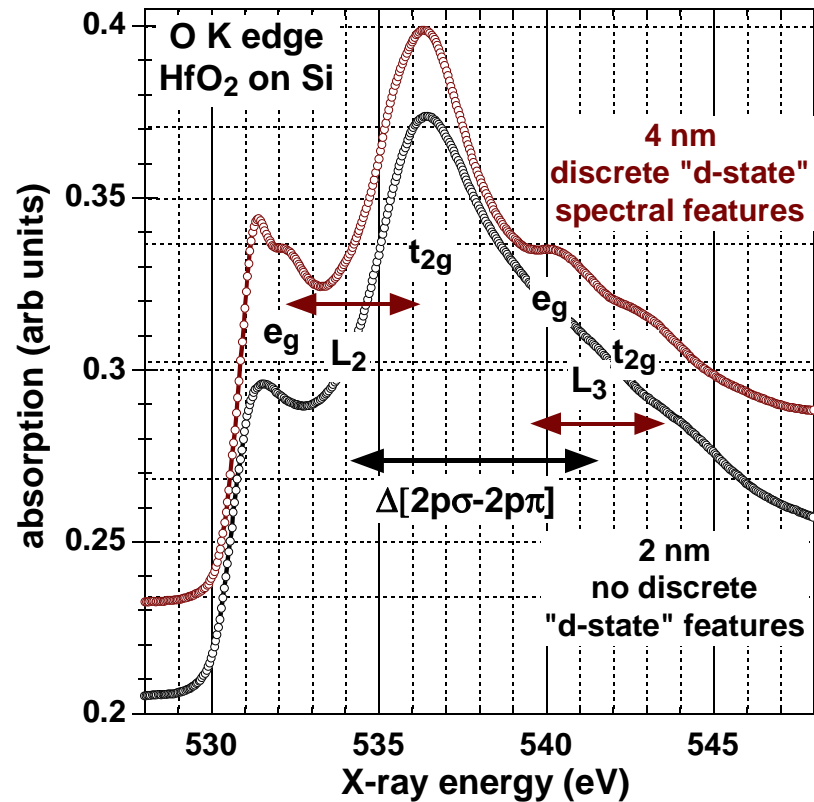
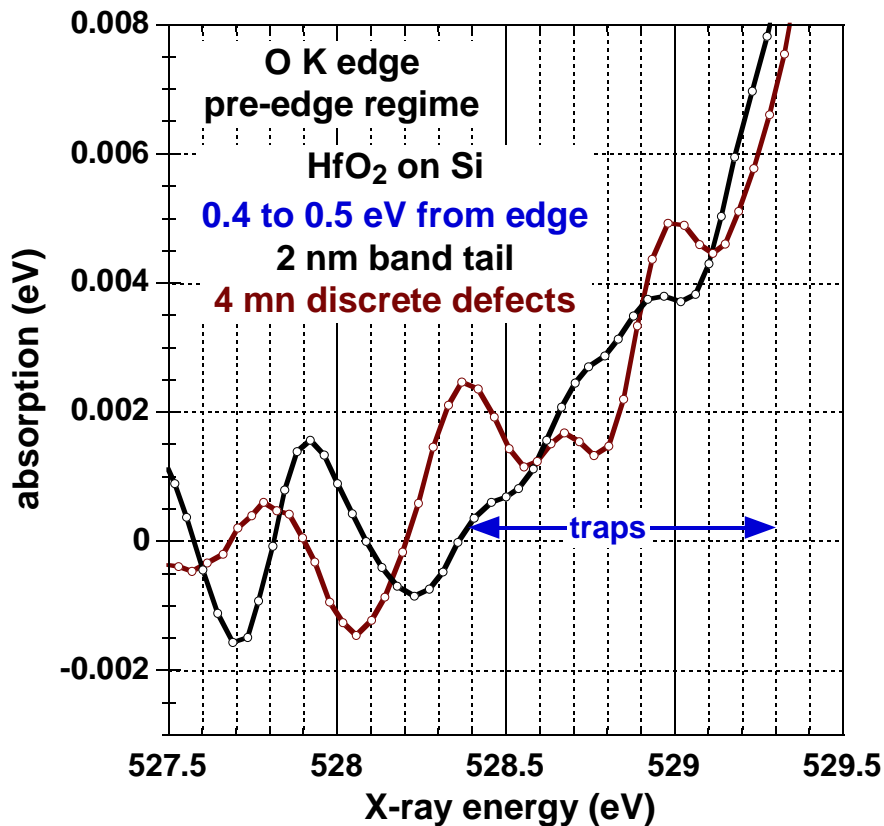
**multiplicity and symmetry of features is determined by symmetry of the primitive unit cell**



**tetragonal  
nano-grains**  
relatively high  
symmetry  
**3 sets of triplets**

**monoclinic  
nano-grains**  
reduced symmetry  
more features  
**mixing of triplet and  
singlet states**

# differences in O-vacancy band edge defects, and "d-state" features in OK edge spectra that correlate with film thickness

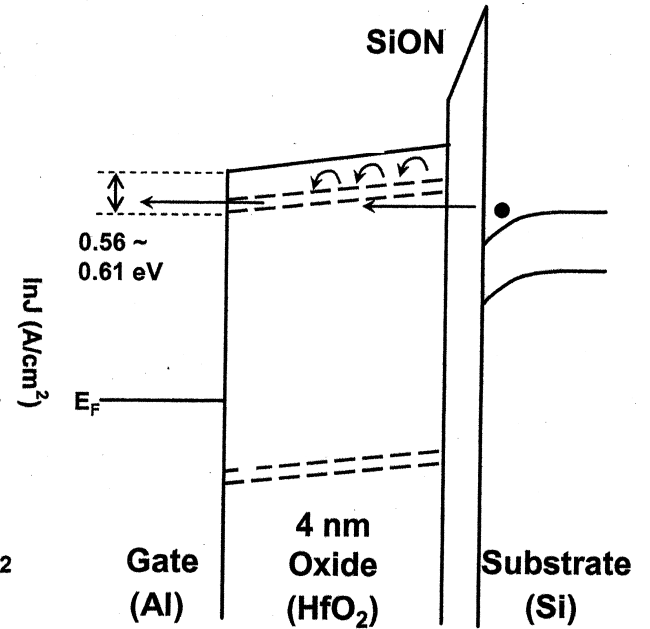
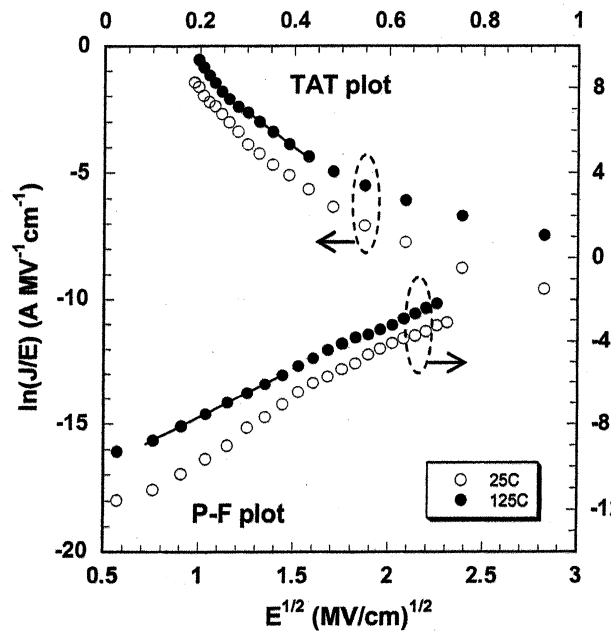
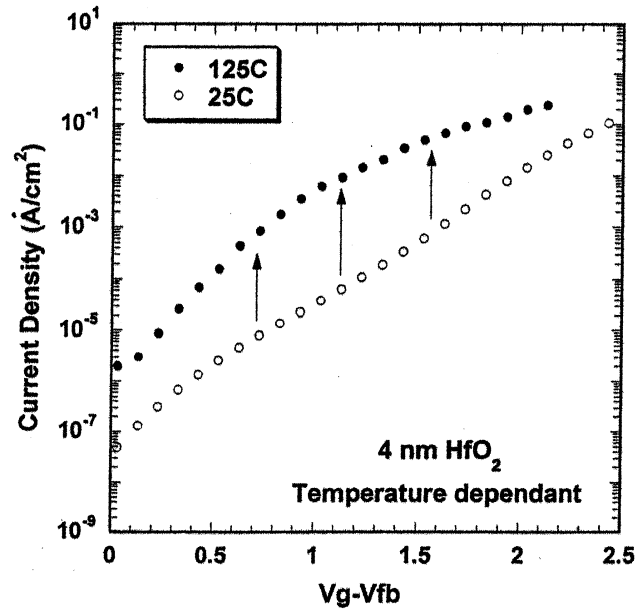
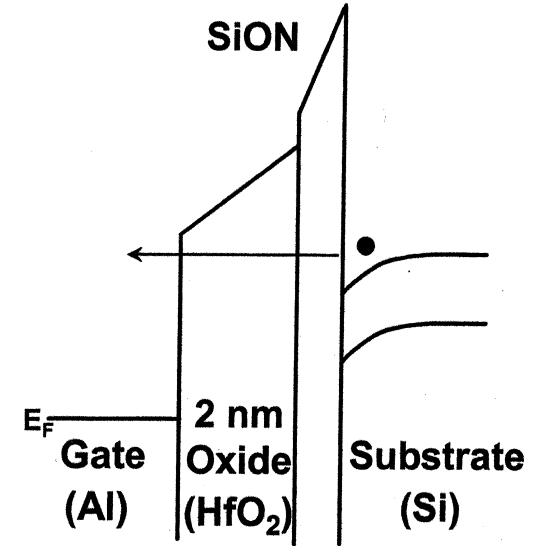
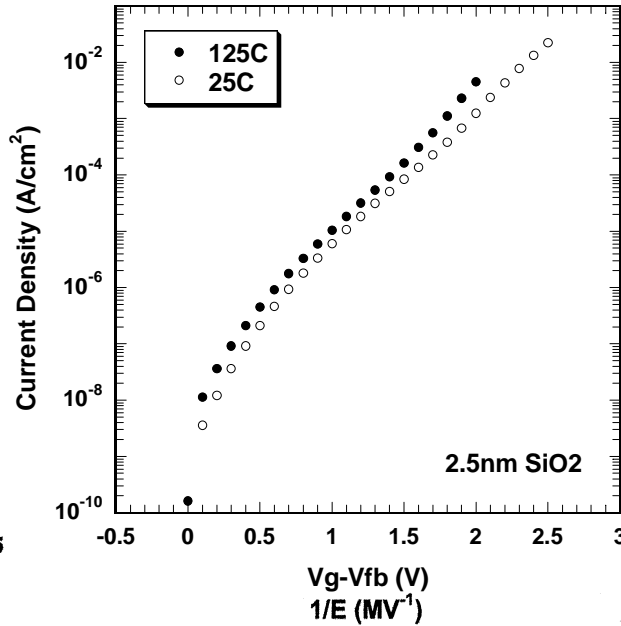
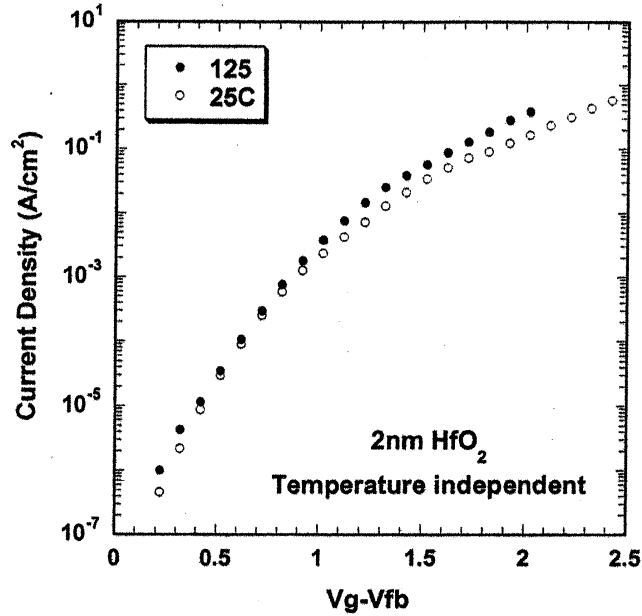


discrete character of defects and "d-state" features depends on unit cell size

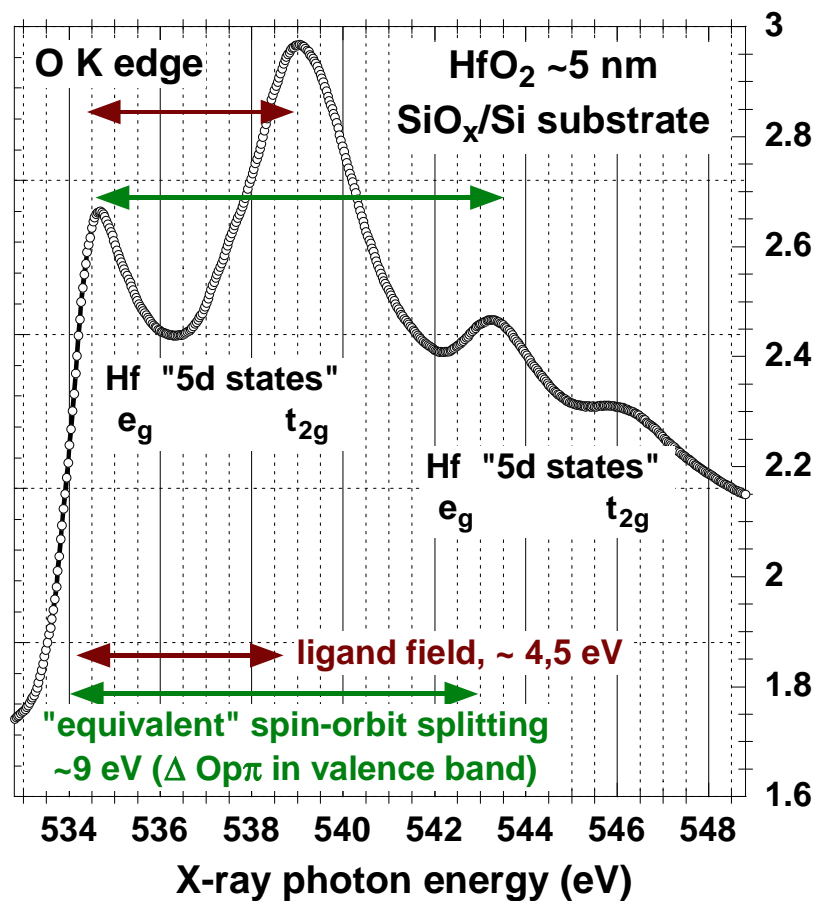
physical thickness of nano-grain films constrains extended unit cell to be less than film thickness

2 nm thick films to less than critical size of ~ 3.0-3.5 nm necessary for cooperative Jahn-Teller distortions

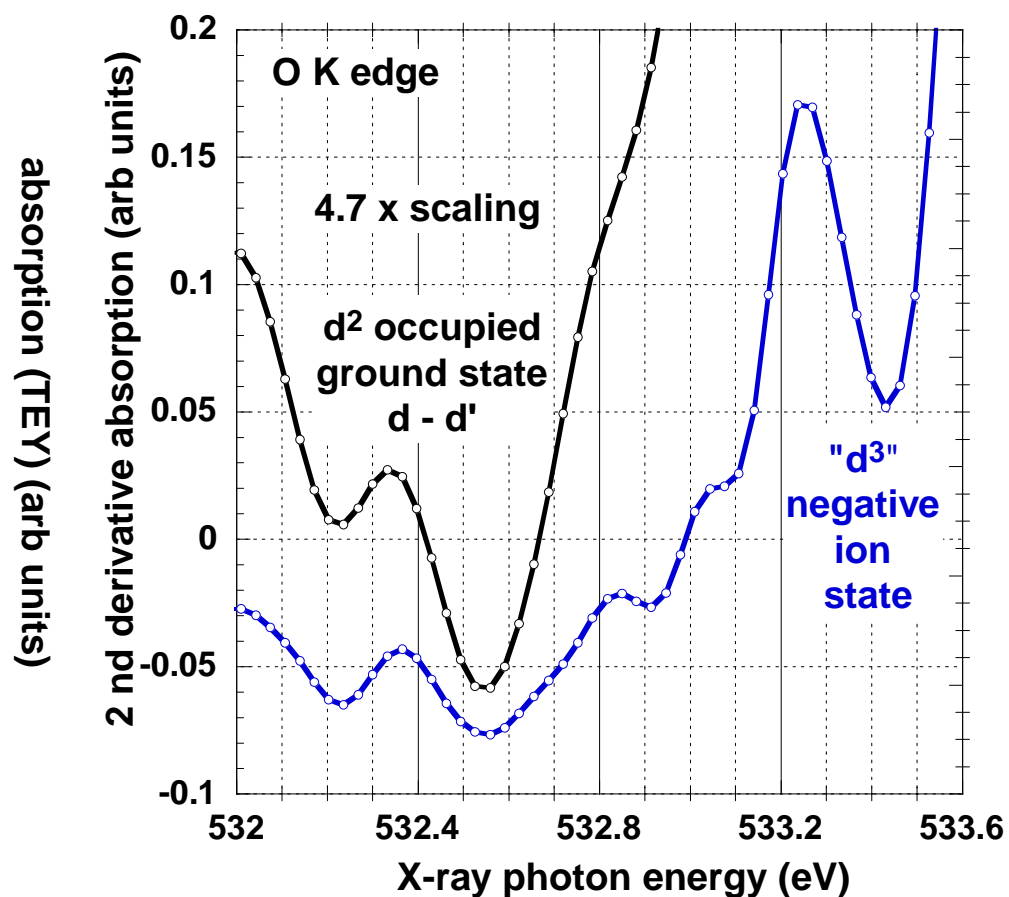
# differences in J-V correlate with differences in conduction band edge and band edge defect states (O-vacancies)



**TEY (total energy yield) SSRL BL 10-1**  
**X-ray absorption spectroscopy - XAS**  
**negative ion states additional shallow traps**  
**closer to conduction band edge**

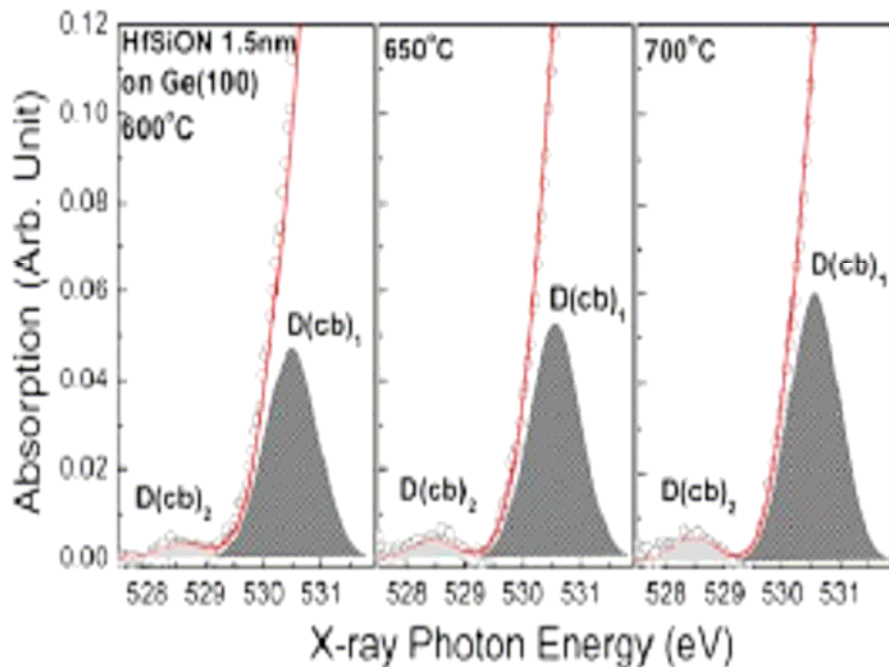
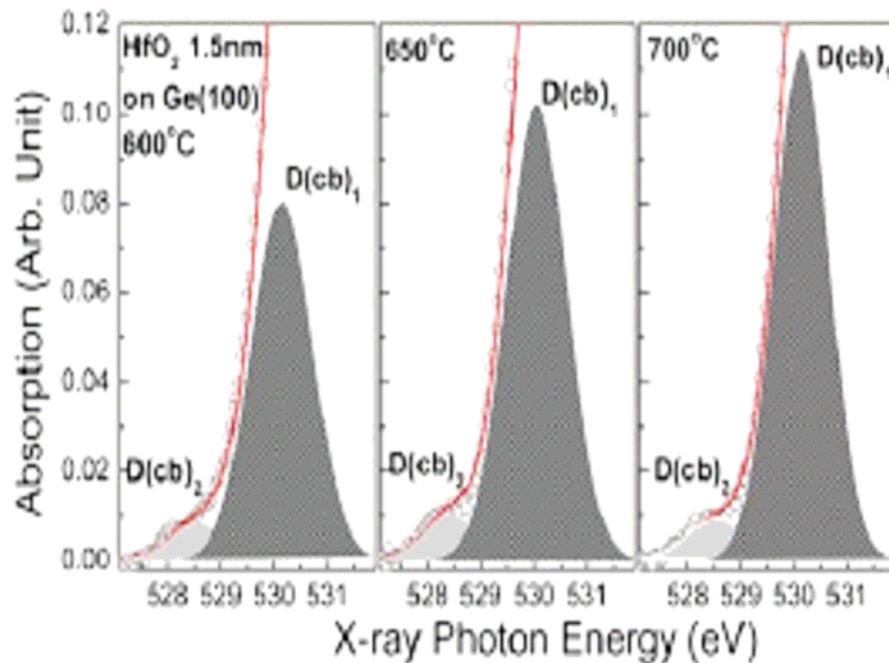


**virtual bound states in  
conduction band**



**pre-edge region virtual bound states**  
**intra-d<sup>2</sup> O-vacancy and**  
**d<sup>3</sup> negative ion state**





**(a) Thermal evolution of the O  $K_1$  spectra for both  $\text{HfO}_2$  and  $\text{HfSiON}$  over a narrow energy region below doubly-degenerate Hf 5d conduction band edge states as a function of PDA temperature.**

**two deconvoluted peaks, labeled  $\text{D}(\text{cb})_1$  and  $\text{D}(\text{cb})_2$ , indicate the defect states determined by Gaussian fits of the XAS O  $K_1$  edge spectra.**

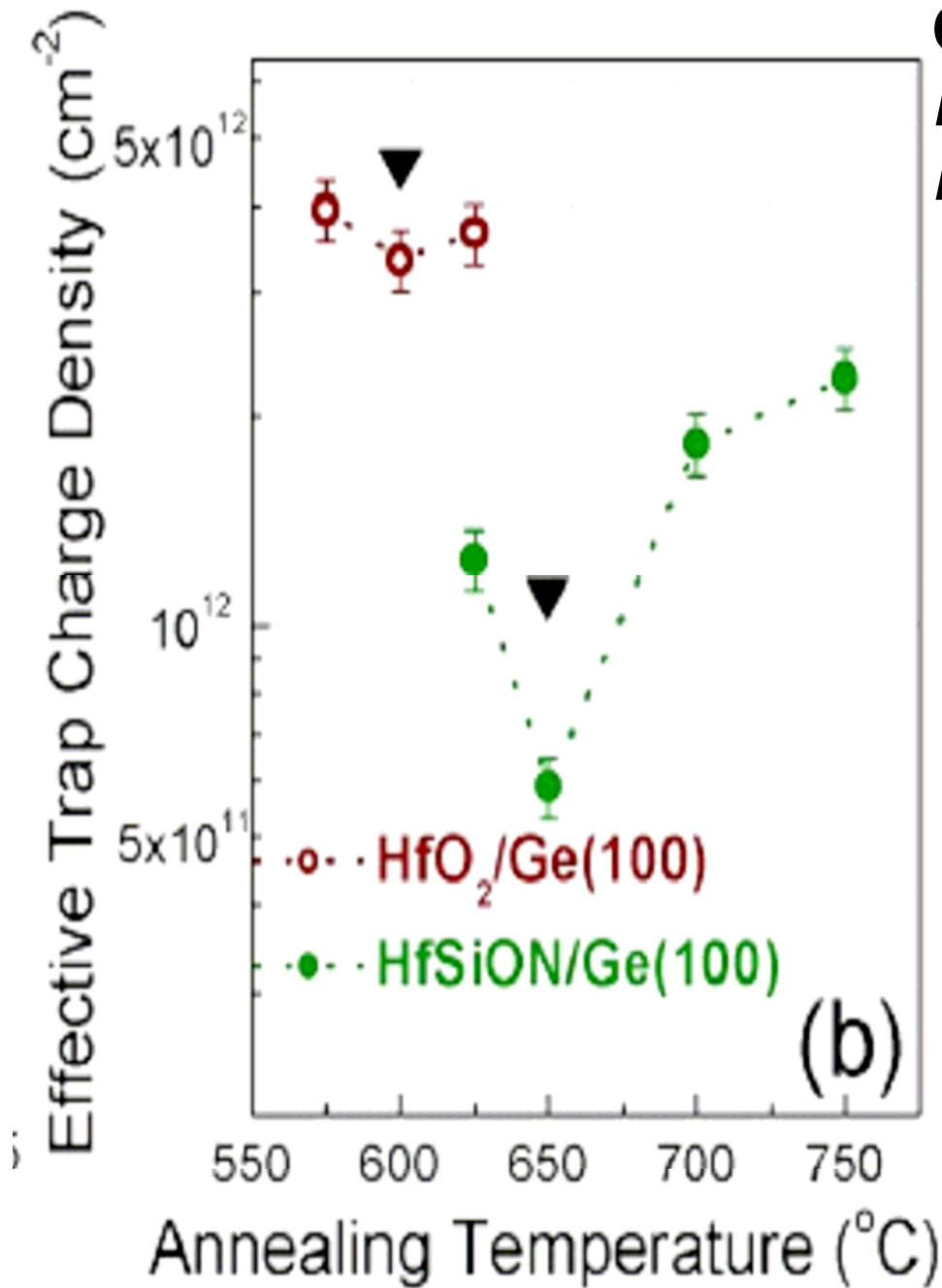
**lower density of defects by spectroscopy for  $\text{HfSiON}$  than  $\text{HfO}_2$  - confirmed by CV**

## Capacitance-voltage (CV) for *n*-MOSCAPS 5nm HfSiON on *n*-type Ge (100)

**lowest trap density same as Saraswat's (Stanford Univ) best**

**HfSiON on Si - X-ray stress similar to SiO<sub>2</sub> no negative fixed charge as in HfO<sub>2</sub>**

**defects lower than in HfO<sub>2</sub> agrees with XAS**



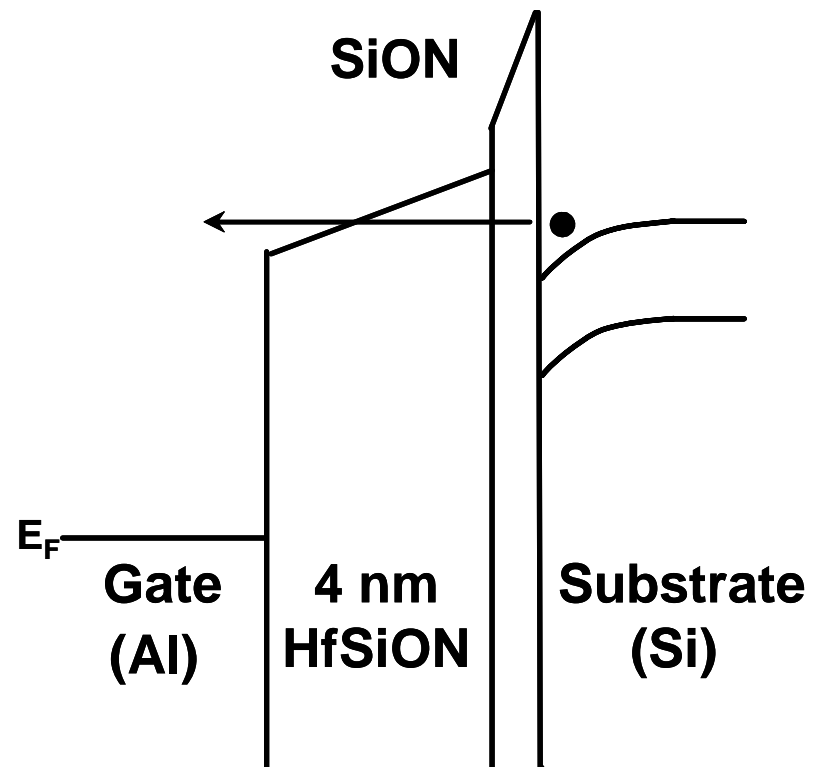
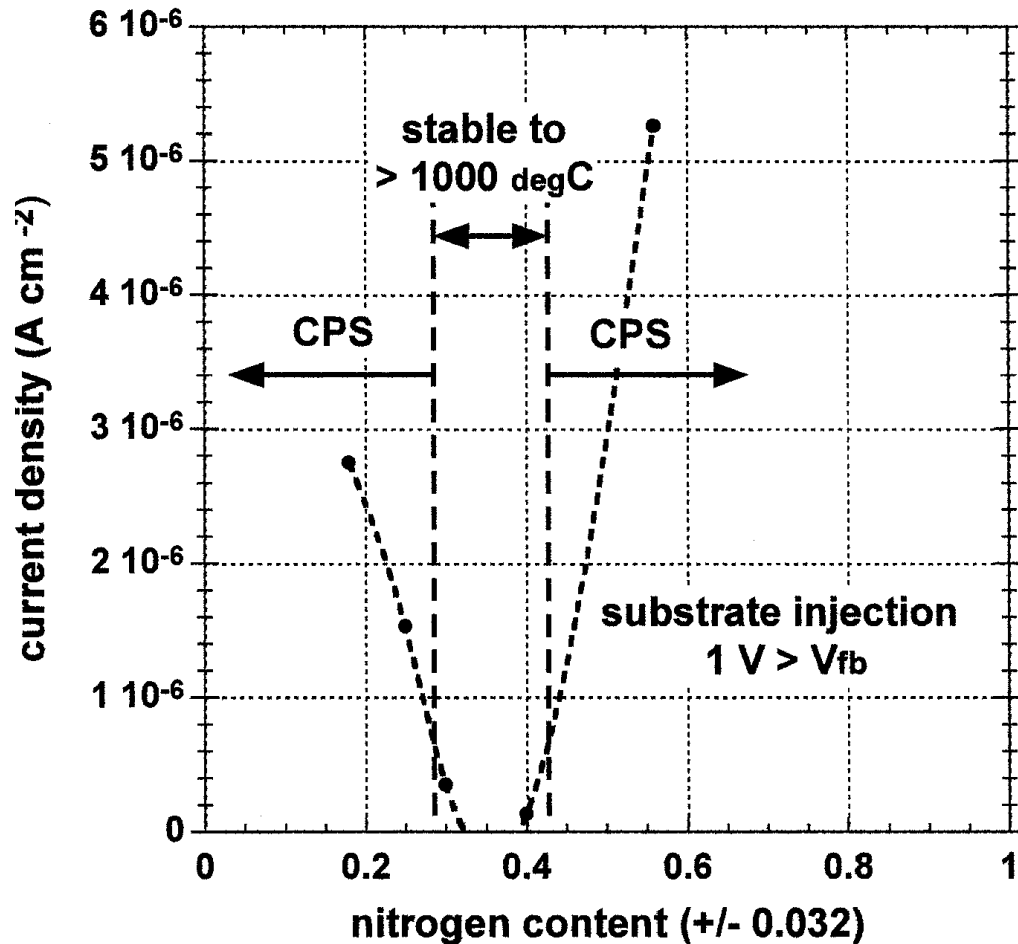




**chemically-controlled composition stable to > 1000°C**

**defect generation - X-ray stress**

**“SiO<sub>2</sub> look alike with a higher dielectric constant ~10-11,  
and a 5000 fold decrease in leakage current !!”**



**HfSiON (Si) requires  
SiON IRT**

\*AFOSR, Vanderbilt MURI Jim Felix, Dan Fleetwood,....

**this review presentation is dedicated to**

**i) Isaac Besuker and Victor Pollinger -- my mentors and friends within the Jahn-Teller theoretical chemistry community,**

**and**

**ii) Rick Garfunkel for his much appreciated critique of a “bad” SRC review presentation a few weeks ago  
this improved talk is my response to his critique**

## **medium range order (SRO) in SiO<sub>2</sub>**

**MRO associated with d-states - non-degenerate A's, B's, doubly degenerate E's, and triply degenerate p-like T<sub>1</sub>'s and d-like T<sub>2</sub>'s**

**this is what makes SiO<sub>2</sub>, and potentially GeO<sub>2</sub>, and their oxynitrides unique with respect to band edge defect states**

**major accomplishment is this fundamental intrinsic connection between many electron theory, synchrotron X-ray spectroscopy and defects**

- dielectrics with band edge A<sub>1g</sub> “s-like” states OK, only positive charge after radiation dosing – SiO<sub>2</sub>**
- HfO<sub>2</sub>, etc. with e<sub>g</sub> or t<sub>2g</sub> - O-vacancy/negative ion electron traps – negative charge after radiation dosing**